Bessel Functions, by Dr Colton Physics 442, Winter 2017

The Bessel functions, $J_{\alpha}(x)$ are a series of functions, that:

- (a) come up often, especially in partial differential equations
- (b) have interesting properties
- (c) are well understood and have been studied for centuries

They are typically only used for positive values of *x*. Here are the first four plotted.



Out[27]= - Graphics -

-0.2

In[28]:= Plot [BesselJ[3, x], {x, 0, 20}]



J₃(x) crosses zero at 6.380, 9.761, 13.015, ...

Out[28]= - Graphics -

Very important facts

- Can be computed via a series formula: $J_{\alpha}(x) = \sum_{k=0}^{\infty} \frac{(-1)^k x^{2k+\alpha}}{k!(k+\alpha)! 2^{2k+\alpha}}$
- Derivatives:
 - $\circ \frac{d}{dx}(J_0(x)) = -J_1(x) \quad \text{(which means the max/min of } J_0 \text{ are at the zeroes of } J_1)$ $\circ \frac{d}{dx}(J_n(x)) = \frac{J_{n-1}(x) J_{n+1}(x)}{2}$ Zeroes: $u_{\alpha m}$ represents the mth zero of $J_{\alpha}(x)$. From the previous page:
- - $u_{01} = 2.405, u_{02} = 5.520, u_{03} = 8.654, \dots$ 0
 - $u_{11} = 3.832, u_{12} = 7.016, u_{13} = 10.173, \dots$
 - 0 etc.

These numbers are available in Mathematica via the BesselJZero function: BesselJZero[0,1] yields a result of 2.4048255576957727686..., etc.

- Differential equations ("Bessel's equation"):
 - $x^2 \frac{d^2 f}{dx^2} + x \frac{df}{dx} + (x^2 \alpha^2)f = 0$ has solution $J_{\alpha}(x)$, or linear combinations 0 $x^2 \frac{d^2 f}{dx^2} + x \frac{d f}{dx} + (u_{\alpha m}^2 x^2 - \alpha^2)f = 0$ has solution $J_{\alpha}(u_{\alpha m} x)$, or linear combinations
 - 0 There are a second set of solutions to Bessel's differential equation, called $Y_a(x)$ (sometimes
- written $N_{\alpha}(x)$) or the "Bessel functions of the second kind"; these diverge at x = 0 and are therefore often not used.
 - Various other linear combinations of $J_{q}(x)$ and $Y_{q}(x)$ are sometimes used; two examples 0 are the "modified Bessel functions" and the "Hankel functions". These are beyond the scope of this course.
 - The so-called "spherical Bessel functions" (and "spherical Hankel functions") are Ο solutions to a different, albeit closely related, differential equation. These are also beyond the scope of this course.
- When u_{am} is added to the argument in the form $u_{am}r$, the functions go to zero at r = 1. Then e.g. $J_0(u_{0m}r)$ represents a series of functions between the origin and a circular boundary of r = 1, having m antinodes. For example $I_0(u_{0,2}x)$ has two antinodes, $I_0(u_{0,3}x)$ has three antinodes, etc.



Orthogonality: $J_0(u_{0m}x)$ and $J_0(u_{0n}x)$ are orthogonal over the interval (0, 1) with respect to a weighting function of *x*:

$$\circ \quad \int_0^1 x J_0(u_{0m}x) J_0(u_{0n}x) dx = \begin{cases} 0 & \text{if } n \neq m \\ \frac{1}{2} (J_1(u_{0m}))^2 & \text{if } n = m \end{cases}$$

Or, more generally for $J_\alpha(x)$:
$$\circ \quad \int_0^1 x J_\alpha(u_{\alpha m}x) J_\alpha(u_{\alpha n}x) dx = \begin{cases} 0 & \text{if } n \neq m \\ \frac{1}{2} (J_{\alpha+1}(u_{\alpha m}))^2 & \text{if } n = m \end{cases}$$

Orthogonality, depicted



 $\label{eq:information} \inf [4] = \mbox{Plot}[\mbox{BesselJ}[0, \mbox{u03 x}], \mbox{x}, \mbox{0, 3}, \mbox{PlotRange} \rightarrow \mbox{All}]$



These are two functions which are orthogonal. They are both J_0 functions, but one has been scaled so that its *second* "zero crossing" occurs at x = 1, and the other has been scalled so that its *third* zero crossing occurs at x = 1.

 $ln[5]:= Plot[x BesselJ[0, u02 x] BesselJ[0, u03 x], \{x, 0, 1\}, Filling \rightarrow Axis]$



In[6]:= Integrate[x BesselJ[0, u02 x] BesselJ[0, u03 x], {x, 0, 1}]

 $\ln[7]:= \operatorname{Plot}[x\operatorname{BesselJ}[0, \operatorname{u03} x] \operatorname{BesselJ}[0, \operatorname{u03} x], \{x, 0, 1\}, \operatorname{Filling} \rightarrow \operatorname{Axis}]$

The two functions are only orthogonal over the range from 0 to 1, and only with a weighting function of x. But within those constraints, you can see that the positive areas cancel out the negative areas.

The integral is zero.

$$\int_0^1 x J_0(u_{0m}x) J_0(u_{0n}x) dx$$

equals 0 if $n \neq m!$

By contrast, $J_0(u_{0m}x)$ is not orthogonal to itself.

In[8]:= Integrate[x BesselJ[0, u03 x] BesselJ[0, u03 x], {x, 0, 1}] // N
Out[8]:= 0.0368432

0.6

0.8

1.0

In[9]:= .5 (BesselJ[1, u03]) ^2 // N
Out[9]:= 0.0368432

0.2

0.4

Out[6]= 0

0.06

0.04 Out[7]=

0.02

 $\int_{0}^{1} x J_{0}(u_{0m}x) J_{0}(u_{0n}x) dx$ equals $\frac{1}{2} (J_{1}(u_{0m}))^{2}$ if n = m!

Sines/Cosines

- Two oscillatory functions: sin(x) and cos(x). Sometimes one of them is not used, due to the symmetry of the problem.
- You determine the value of sin(x) or cos(x) for arbitrary x by using a calculator or computer program.

3.
$$\sin(x) = \sum_{k=0}^{\infty} \frac{(-1)^k x^{2k+1}}{(2k+1)!}$$

 $\cos(x) = \sum_{k=0}^{\infty} \frac{(-1)^k x^{2k}}{(2k)!}$

Consider just sin(x):

- 4. The zeroes of sin(x) are at $x = \pi, 2\pi, 3\pi$, etc. $x = "m\pi"$ is the *m*th zero
- 5. Using the zeroes in the argument, $sin(m\pi x)$ has *m* antinodes in the interval from 0 to 1. And at x = 1 (the boundary), $sin(m\pi x) = 0$ for all *m*.
- 6. The differential equation satisfied by f = sin(x) is f'' + f = 0.

The differential equation satisfied by $f = \sin(m\pi x)$ is $f'' + (m\pi)^2 f = 0$.

7. $\sin(n\pi x)$ is orthogonal to $\sin(m\pi x)$ on the interval (0,1):

Two oscillatory functions for each α : $J_{\alpha}(x)$ and $Y_{\alpha}(x)$. Typically Y_{α} is not used because it's infinite at the origin.

You determine the value of $J_{\alpha}(x)$ for arbitrary x by using a calculator or computer program.

$$J_{\alpha}(x) = \sum_{k=0}^{\infty} \frac{(-1)^{k} x^{2k+\alpha}}{k!(k+\alpha)! 2^{2k+\alpha}}$$

<u>Consider just $J_{\alpha}(x)$ for one α , say $\alpha = 0$: (similar things hold true for all α 's) The zeroes of $J_0(x)$ are at $x \approx 2.405, 5.520, 8.654$, etc. $x = "u_{0m}"$ is the m^{th} zero Using the zeroes in the argument, $J_0(u_{0m}r)$ has m</u>

antinodes in the circular region $r \le 1$. And at r = 1 (the boundary), $J_{\theta}(u_{0m}x) = 0$ for all *m*.

The differential equation satisfied by $f = J_0(x)$ is $x^2 f'' + x f' + (x^2 - 0^2) f = 0$.

The differential equation satisfied by $f = J_0(u_{0m}x)$ is $x^2 f'' + x f' + (u_{0m}^2 x^2 - 0^2) f = 0.$

 $0^2 \rightarrow \alpha^2$ for other α 's

 $J_0(u_{0n}x)$ is orthogonal to $J_0(u_{0m}x)$ on the interval (0,1), with respect to a weighting function of *x*:

$$\int_{0}^{1} \sin(n\pi x) \sin(m\pi x) dx = \begin{cases} 0, & \text{if } n \neq m \\ \frac{1}{2}, & \text{if } n = m \end{cases}$$

$$\int_{0}^{1} x J_{0}(u_{0n}x) J_{0}(u_{0m}x) dx = \begin{cases} 0, if \ n \neq m \\ \frac{1}{2} (J_{1}(u_{0m}))^{2}, if \ n = m \end{cases}$$

Similar orthogonality holds for other values of α .

Additionally, the Bessel functions are related to sines/cosines through this integral formula:

$$J_{\alpha}(x) = \frac{1}{\pi} \int_{0}^{\pi} \cos(\alpha \theta - x \sin \theta) d\theta$$

Quote from Mary Boas, in *Mathematical Methods in the Physical Sciences*: "In fact, if you had first learned about sin(nx) and cos(nx) as power series solutions of $y'' = -n^2y$, instead of in elementary trigonometry, you would not feel that Bessel functions were appreciably more difficult or strange than trigonometric functions. Like sines and cosines, Bessel functions are solutions of a differential equation; they are tabulated and their graphs can be drawn; they can be represented as a series; and a large number of formulas about them are known."