

Lecture 9: Mon, 28 Jan 2008

Reading quizzes: no talking, no looking in your books/notes

Q1. Brewster's angle occurs when θ_i and θ_t

- a. Are equal
- b. Sum to 90°
- c. Sum to 45°
- d. Sum to 60°

Q2. During total internal reflection, the electric field in the air just outside the glass is called the

- a. reduction wave
- b. florescent wave
- c. evanescent wave
- d. vanishing wave
- e. nothing...there can be no E-field there.

Q3. Total internal reflection happens for:

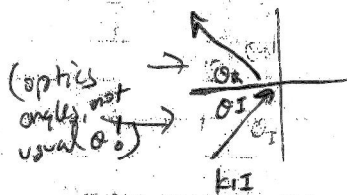
- a. s-polarized light
- b. p-polarized light
- c. both polarizations

July 9

Also $k_{Iy} = k_{Ry}$

$k_{Iy} = k_{Ry} \rightarrow$ incident and reflected have same y-component!

$\rightarrow$ same magnitude, same y-component



only possible if

$\theta_R = \theta_I$ 2nd Law!

"Law of Reflection"

Also $k_{Iy} = k_{Ty}$

$v = \frac{\omega}{k}$

$k = \frac{\omega}{v}$

$k = \frac{\omega n}{c}$

$k_1 \sin \theta_1 = k_2 \sin \theta_2$

$\frac{\omega}{k} = v = \frac{c}{n}$

$\omega \frac{n_1}{c} \sin \theta_1 = \omega \frac{n_2}{c} \sin \theta_2$

$n_1 \sin \theta_1 = n_2 \sin \theta_2$

3rd Law

"Law of Refraction", "Snell's Law"

(also critical angle $\rightarrow$ can't get any transmission)

3 Laws of Geometrical Optics!

Gives deeper explanation: a beam of light bends because its fields propagate at a different speed (like driving on shoulder)

Explains focusing, magnification, ray tracing, etc.

$\rightarrow$ only used most basic boundary condition, not really any advanced Maxwell Eq

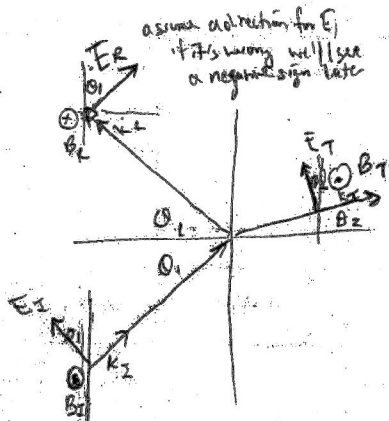
So these are true for basically any waves!

$\rightarrow$ you're starting to realize why it was so hard to accept Einstein's photon idea!

Back to Basic Problem

Now let's assume a particular polarization: p-polarization

$\vec{E} \parallel$ to plane of incidence



BC $E_i = E_r + E_t$ $\tilde{E}_i \cos \theta_i + \tilde{E}_r \cos \theta_r = \tilde{E}_t \cos \theta_t$ (lost all exponential factors because they're equal)

BC $B_i = B_r + B_t$ $\frac{1}{\mu_1 v_1} \tilde{E}_i = \frac{1}{\mu_1 v_1} \tilde{E}_r + \frac{1}{\mu_2 v_2} \tilde{E}_t$
 $\tilde{E}$ in opposite direction

let $\beta = \frac{\mu_1 v_1}{\mu_2 v_2}$ as before $\frac{n_2}{n_1}$

then $\tilde{E}_i - \tilde{E}_r = \beta \tilde{E}_t$

let $\alpha = \frac{\cos \theta_t}{\cos \theta_r}$ then $\tilde{E}_i + \tilde{E}_r = \alpha \tilde{E}_t$

Solve simult. eqns with a little algebra

$$\tilde{E}_r = \left(\frac{\alpha - \beta}{\alpha + \beta} \right) \tilde{E}_i$$

$$\tilde{E}_t = \frac{2}{\alpha + \beta} \tilde{E}_i$$

Fresnel eqns! 3.15 + 3.16
 [same as P+W eqns, multiply num + denom by $\cos \theta_r$ to get right most form]

Notice: a 180° phase shift in reflected beam if $\beta > \alpha$

if $\beta = \alpha$ special case, "Brewster's Angle" $\rightarrow \tilde{E}_r = 0$! no reflection

$$\frac{\cos \theta_t}{\cos \theta_r} = \frac{n_2}{n_1}$$

$$\sin \theta_t = \frac{n_1}{n_2} \sin \theta_r$$

$$\frac{\sqrt{1 - \sin^2 \theta_t}}{\cos \theta_r} = \frac{n_2}{n_1} = \cos \theta_r$$

$$\sqrt{1 - \frac{n_1^2}{n_2^2} \sin^2 \theta_r} = \frac{n_2}{n_1} \cos \theta_r$$

$$1 - \frac{n_1^2}{n_2^2} \sin^2 \theta_r = \frac{n_2^2}{n_1^2} \cos^2 \theta_r$$

FW 10/15/14

Can skip unless needed

$$1 - \frac{1}{\beta^2} \sin^2 \theta_1 = \beta^2 \cos^2 \theta_1 \quad (\times n^2)$$

$$\beta^2 - \sin^2 \theta_1 = \beta^4 (1 - \sin^2 \theta_1)$$

$$\beta^4 \sin^2 \theta_1 - \sin^2 \theta_1 = \beta^4 - \beta^2$$

$$\sin^2 \theta_1 (\beta^4 - 1) = \beta^4 - \beta^2$$

$$\sin^2 \theta_1 = \frac{\beta^2 (\beta^2 - 1)}{\beta^4 - 1} = \frac{\beta^2 (\beta^2 - 1)}{(\beta^2 + 1)(\beta^2 - 1)}$$

$$\sin^2 \theta_1 = \frac{\beta^2}{\beta^2 + 1}$$

$$\text{consider } \cos^2 \theta_1 = 1 - \sin^2 \theta_1 = 1 - \frac{\beta^2}{\beta^2 + 1} = \frac{\beta^2 + 1 - \beta^2}{\beta^2 + 1} = \frac{1}{\beta^2 + 1}$$

$$\text{so } \tan^2 \theta_1 = \frac{\frac{\beta^2}{\beta^2 + 1}}{\frac{1}{\beta^2 + 1}} = \beta^2$$

$$\tan \theta_B = \beta = \frac{n_2}{n_1}$$

Brewster angle condition!

only a 180° phase shift for certain angles ($\theta < \theta_B$)

Gravitational Boost Figs

→ phase shift in Refl for $\theta < \theta_B$ (if normal???)

P7391 and 392

• For $\theta_I = 90^\circ$ (grazing), $\alpha = \frac{\cos \theta}{0} = \infty$

→ all reflected, none transmitted!

• For $\theta_I = 0^\circ$ (normal) → $\alpha = 0$, recover previous Fresnel Eqs

Reflected intensities: $I = \frac{1}{2} \epsilon_1 v_1 E_{01}^2 \rightarrow$ not good enough, we need $\vec{S} \cdot \hat{z}$

$$R = \frac{E_{0R}^2}{E_{0I}^2} = \left(\frac{\alpha - \beta}{\alpha + \beta} \right)^2 \quad \epsilon's, v's, \theta's \text{ equal}$$

$$T = \underbrace{\frac{n_2}{n_1}}_{\beta} \underbrace{\left(\frac{E_{0T}}{E_{0I}} \right)^2}_{\left(\frac{2}{\alpha + \beta} \right)^2} \underbrace{\cos \theta_T}_{\alpha} = \frac{4\alpha\beta}{(\alpha + \beta)^2}$$

Power
unitless of transmission

$$\frac{1}{v_1 \cos \theta} = \alpha$$

From Griffiths

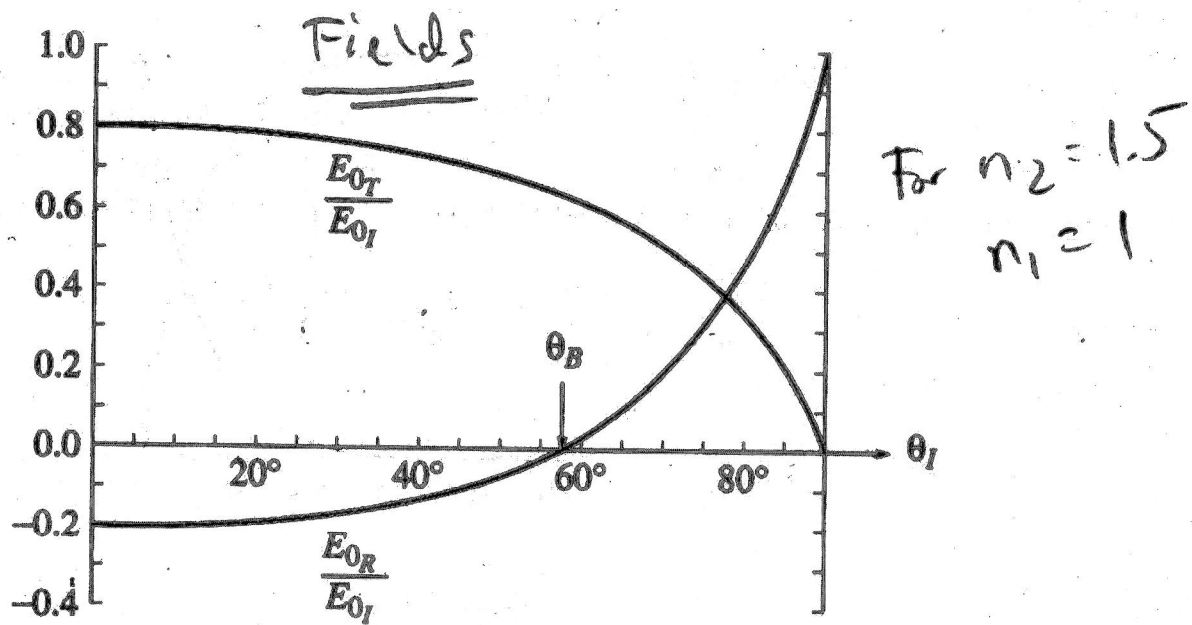


Figure 9.16

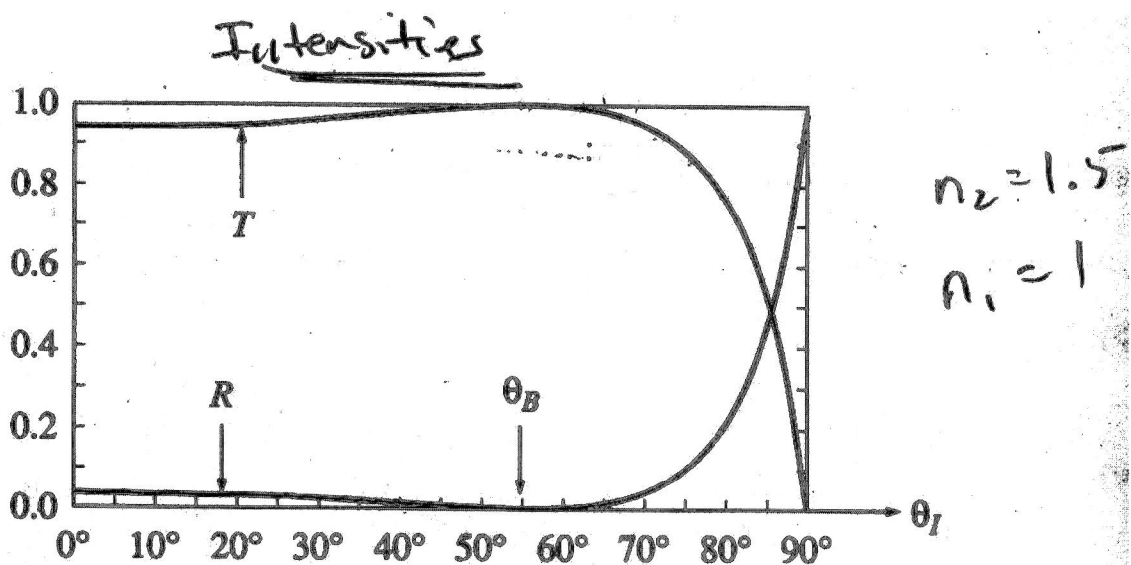
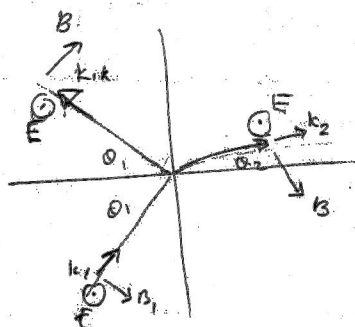


Figure 9.17

S-polarization



$$\tilde{E}_{0R} = \frac{1 - \alpha\beta}{1 + \alpha\beta} \tilde{E}_{0I}$$

$$\tilde{E}_{0T} = \frac{2}{1 + \alpha\beta} \tilde{E}_{0I}$$

Result:

$$R = \left(\frac{1 - \alpha\beta}{1 + \alpha\beta} \right)^2$$

$$T = \frac{4\alpha\beta}{(1 + \alpha\beta)^2}$$

Can do analysis for Hs!

Phase shift in reflected if $n_2 > n_1$
 if $n_2 < n_1$

→ But no Brewster's angle (always get some reflection)

(unless weird μ 's)

Total Internal Reflection

Bring back + Q P ② $n_1 > n_2$



Review

When Snell's law predicts $\theta_2 = 90^\circ$
it can't bend anymore

$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$

$$\text{critical } \sin \theta_1 = \frac{n_2}{n_1}$$

$$\theta_1 \geq \sin^{-1}\left(\frac{n_2}{n_1}\right) \text{ TIR occurs}$$

Proof that it's total reflection

For p polar. $r = \frac{\tilde{E}_{0r}}{\tilde{E}_{0i}} = \frac{\alpha - \beta}{\alpha + \beta}$

$$\alpha = \frac{\cos \theta_2}{\cos \theta_1}$$

$$\beta = \frac{n_2}{n_1}$$

α is now imaginary! But that's OK with equations

$$\begin{aligned} \text{proof: } \alpha &= \frac{\sqrt{1 - \sin^2 \theta_2}}{\cos \theta_1} = \frac{\sqrt{1 - \left(\frac{n_1}{n_2}\right)^2 \sin^2 \theta_1}}{\cos \theta_1} \rightarrow \text{larger than 1} \\ &= \frac{\sqrt{\text{negative}}}{\text{positive}} \\ &= \text{Imaginary!} \\ &= ix \end{aligned}$$

What does that do to r ?

$$r = \frac{ix - \beta}{ix + \beta}$$

What is magnitude of r ? $|r| = \sqrt{r r^*}$

$$\begin{aligned} &= \sqrt{\left(\frac{ix - \beta}{ix + \beta}\right) \left(\frac{-ix - \beta}{-ix + \beta}\right)} \\ &= \sqrt{(-1)(-1)} \\ &= 1 \checkmark \end{aligned}$$

evanescent field
→ long eqn
p 64

All is reflected!

Optics
animation!
animation 5

Ref from cond. charge n now can be complex $n = n + ik$

→ not only is α imaginary (complex?)
but β is also!

but aside from that things are same

$r \neq 1$ though

Lecture 10: Wed, 30 Jan 2008

Reading quizzes: no talking, no looking in your books/notes

Q1. In circularly polarized light, the two components of the E-field differ by a phase of

- a. 0
- b. $\pi/3$
- ☒ c. $\pi/2$
- d. $2\pi/3$
- e. π

Q2. The mathematical objects that describe polarization states are

_____ vectors

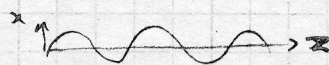
- a. Fred
- ☒ b. Jones
- c. Newton
- d. Maxwell
- e. Smith

Q3. In the above vectors, the ones that use complex numbers represent

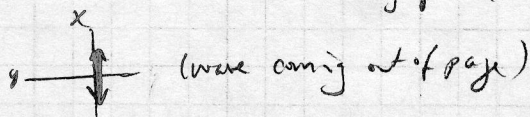
- a. linear polarized light
- ☒ b. elliptical light in general
- c. circular only

Polarization

We've been studying things like $\vec{E} = \vec{E}_0 e^{i(kz - \omega t)}$
 = (say) $E_0 \hat{x} \cos(kz - \omega t)$



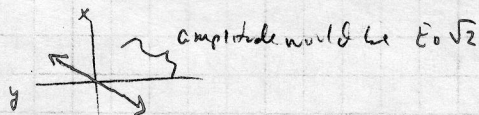
This is a "linearly polarized" wave, because the electric field goes back + forth in a line in the x-y plane



What happens when we ^{add} two waves together?

Example: $E_0 \hat{x} \cos(kz - \omega t) + E_0 \hat{y} \cos(kz - \omega t)$

$$= E_0 \sqrt{2} \left(\frac{\hat{x} + \hat{y}}{2} \right) \cos(kz - \omega t)$$



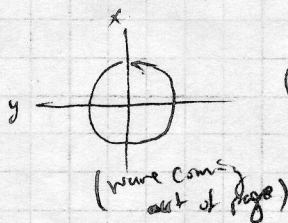
Still linear polarization!

Add a phase shift:

$$E_0 \hat{x} \cos(kz - \omega t) + E_0 \hat{y} \sin(kz - \omega t)$$

$\cos(kz - \omega t - 90^\circ)$

Circular pol! See picture!
 ("amplitude" = E_0 , not $E_0 \sqrt{2}$)

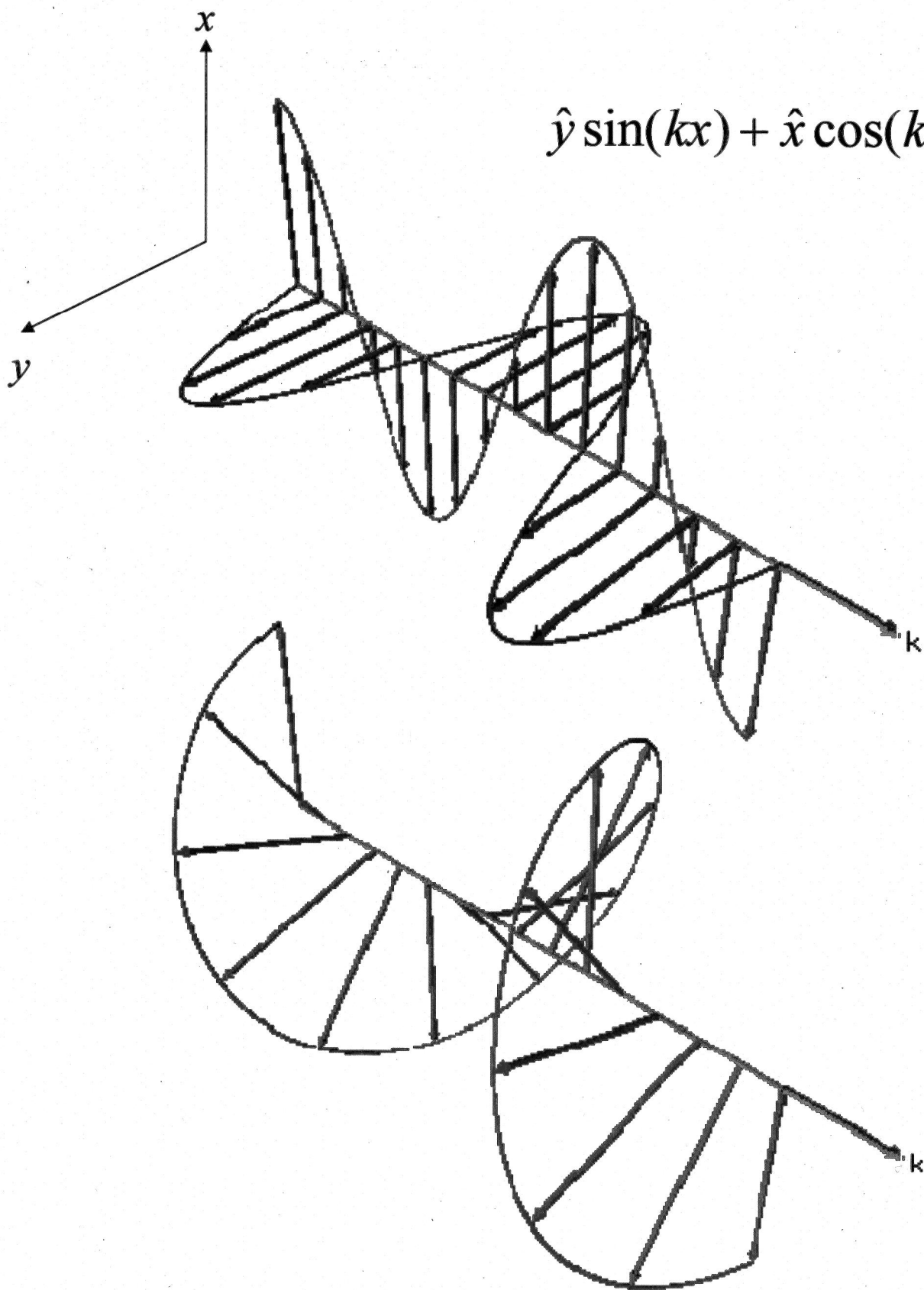


(right circular, to be precise)

stop time, look at screw in E field
 same as regular screw → right handed
 opposite → left handed

also: right handed: $\vec{E}$ rotates to right
 as seen from observer

$$\hat{y} \sin(kx) + \hat{x} \cos(kx)$$



writing circ. pol w/ complex numbers.

$$\begin{aligned}\vec{E}_0 &= E_0 \hat{x} \cos(kz - \omega t) + E_0 \hat{y} \cos(kz - \omega t - 90^\circ) \\ &= E_0 \hat{x} e^{i(kz - \omega t)} + E_0 \hat{y} e^{i(kz - \omega t - 90^\circ)} \\ &= [E_0 \hat{x} + E_0 \hat{y} e^{i(-90^\circ)}] e^{i(kz - \omega t)}\end{aligned}$$

$$E_0 \begin{pmatrix} 1 \\ e^{i(-90^\circ)} \end{pmatrix} e^{i(kz - \omega t)}$$

$\swarrow$
 $e^{i(-90^\circ)} = -i$

The "Jones vector"!

[conversion: vector has magnitude 1, so add $\sqrt{2}$]

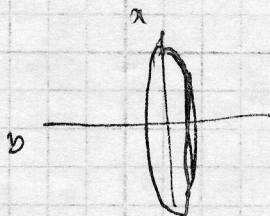
What if the y component not as strong as x-component?

$$E_0 = (E_0 \hat{x} + .2 E_0 \hat{y} e^{i(-90^\circ)}) e^{i(kz - \omega t)}$$

$$J_{\text{res}} = \frac{1}{1.0198} \begin{pmatrix} 1 \\ .2 e^{i(-90^\circ)} \end{pmatrix}$$

$$\frac{1}{1.0198} \begin{pmatrix} 1 \\ .2 \end{pmatrix}$$

Answer #3



Elliptically polarized

$$\text{ellipticity} = \frac{E_{\min}}{E_{\max}}$$

(like "eccentricity" of ellipse, but not quite the same mathematically)

Figure

Summary table 4.01

$$\begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad x$$

$$\begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad y$$

$$\begin{pmatrix} \cos \alpha \\ \sin \alpha \end{pmatrix} \quad \angle \alpha \text{ from } x$$

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix} \quad \text{RCP}$$

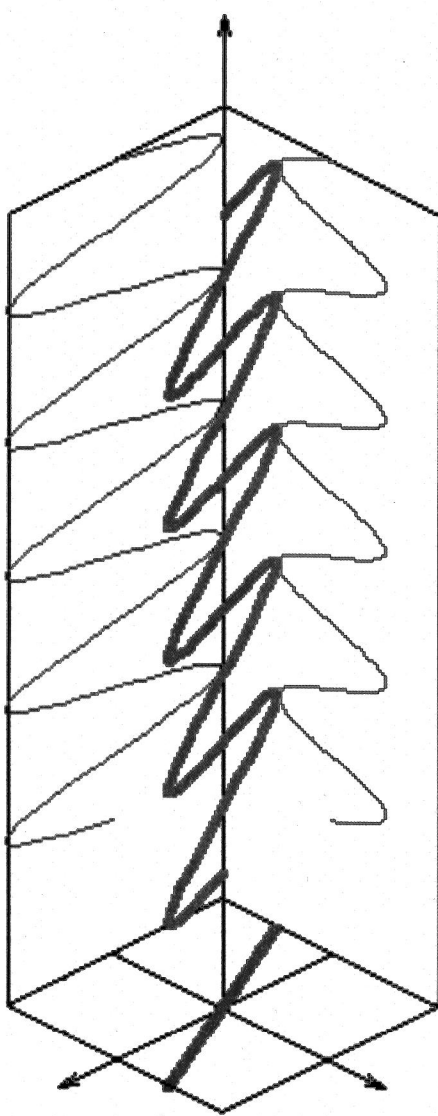
$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix} \quad \text{LCP}$$

$$\frac{1}{\sqrt{A^2 + B^2}} \begin{pmatrix} A \\ B e^{i\phi} \end{pmatrix} \quad \text{RCP-like elliptic, long axis oriented at } \alpha \text{ or } y \text{ axis}$$

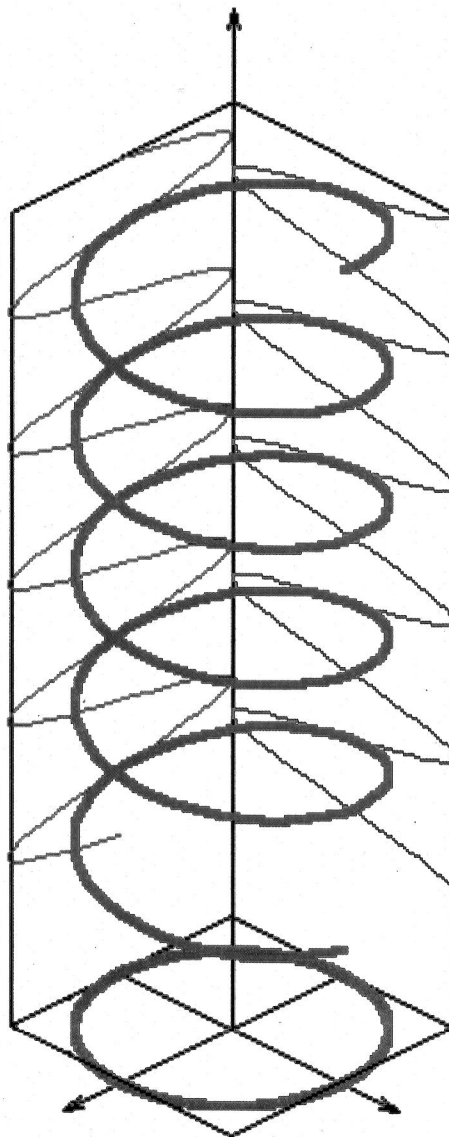
$$\frac{1}{\sqrt{A^2 + B^2}} \begin{pmatrix} A \\ B e^{i\phi} \end{pmatrix} \quad \text{elliptic at inclination } \alpha \text{ (angle } \alpha \text{ given in text)}$$

Elliptically polarized light

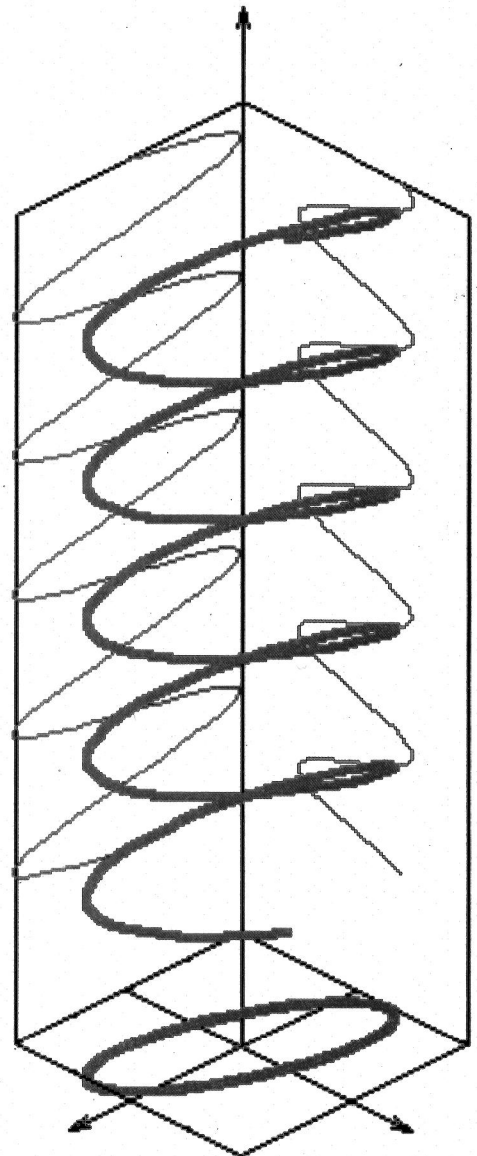
E still rotates at a given point in space. But E is strongest when it points along a certain line, and weakest when it points 90 degrees from that line.



linear



circular



elliptical

Jones vectors summary (similar to P&W table 4.1)

$$\begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

linearly polarized in x-direction

$$\begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

linearly polarized in y-direction

$$\begin{pmatrix} \cos \alpha \\ \sin \alpha \end{pmatrix}$$

linearly polarized at angle α from x-axis

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ e^{i(-90^\circ)} \end{pmatrix}$$

RCP

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ e^{i(+90^\circ)} \end{pmatrix}$$

LCP

$$\begin{pmatrix} A \\ Be^{i(-90^\circ)} \end{pmatrix}$$

RCP-like elliptical, with ellipse oriented

along x- and y-axes ($A^2 + B^2$ must = 1)

$$\begin{pmatrix} A \\ Be^{i\delta} \end{pmatrix}$$

elliptical oriented at angle α from x- or y-axis (long equation for α as a function of δ given in textbook; $A^2 + B^2$ must = 1)