

Sec 10.5 Fraunhofer Approx.

$$e^{i\frac{k}{2d}(x'^2+y'^2)} \approx 1$$

True if $d \rightarrow \infty$ "far field"

What about $e^{-i\frac{k}{d}(x'x''+y'y'')}$ term? Not also = 1?

\rightarrow Not necessarily: x' and y' limited by size of aperture
 x and y have no such limit.

Result

$$E \approx \frac{-i}{\lambda d} e^{ikd} e^{i\frac{k}{2d}(x^2+y^2)} \iint_{\text{aperture}} E(x'',y'') e^{-i\frac{k}{d}(x'x''+y'y'')} dx'' dy''$$

units of wavevector (m⁻¹), but Not a wavevector

$$k_x = \frac{kx}{d}$$

$$k_y = \frac{ky}{d}$$

$k \approx \frac{2\pi}{\lambda}$
normal wavevector

$$\iint E(x'') e^{-i(k_x x'' + k_y y'')} dx'' dy''$$

2D Fourier Transform of aperture function! (w/o 2π)

(or inverse FT depending on convention)

x'' like t
 k_x like ω

We mainly want intensity, $I \propto |E|^2$. Stuff in front of integral except $\frac{1}{\lambda d}$ goes away

Also Consider $\left| \int f(x) e^{ikx} dx \right|^2$ vs $\left| \int f(x) e^{-ikx} dx \right|^2$

overall brightness/darkness Not pattern

Therefore FT vs IFT doesn't affect intensity. They are equal! Because $\| \cdot \|^2 = [\cdot] \times [\cdot]^*$

And since $f(x)$ is real, each $[\cdot]$ is CC of the other

important! \star

$$I = I_0 \times |2D \text{ F.T. of aperture function}|^2$$

\uparrow can even lump $\frac{1}{\lambda d}$ factors into I_0

often aperture function $\begin{cases} 1 & \text{inside aperture} \\ 0 & \text{outside aperture} \end{cases}$
 can be otherwise, e.g. filter over part of aperture

2D Fourier Transforms


often our aperture function will be separable
 $E(x,y) = E(x') E(y')$

then

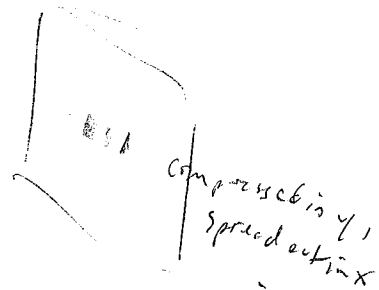
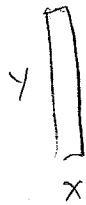
$$\iint E(x,y) e^{ik_x x'} e^{ik_y y'} dx' dy'$$

$$= \int E(x') e^{ik_x x'} dx' \int E(y') e^{ik_y y'} dy'$$

$$= FT(x \text{ func}) \cdot FT(y \text{ func})$$

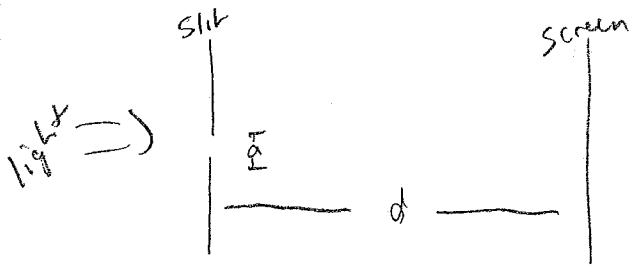
Example: rectangular slit 

If one dimension is long then ~~the~~ structure dominated
 by other dimension



See Mathematics
 plots

Application: Single Slit



Aperture function AF

$$FT(\Pi_a) = \int_{-a/2}^{a/2} e^{ikx} dx = \frac{1}{ik} e^{ikx} \Big|_{-a/2}^{a/2} = \frac{1}{ik} (e^{ika/2} - e^{-ika/2}) = \frac{2 \sin(ka/2)}{k}$$

$= \frac{a \text{sinc}(ka/2)}{k}$

What is diffraction pattern of a single slit?

If it's long + thin, we can neglect the Y-dimension

$$E \propto \hat{F}(AF)$$

~~This AF is the "rectangular pulse"~~

~~studied early, scaled to width a, so~~

$$E = E_0 \text{sinc}\left(\frac{kx a}{2}\right)$$

X, Y = coordinates on screen

Using defn of $k_x = \frac{kX}{d}$

$$I(X) = I_0 \text{sinc}^2\left(\frac{k a X}{2d}\right)$$

→ this is the pattern you see!

If it's rectangular, but not long enough to neglect the Y-dimension, you just get another sinc for Y.

$$I(x) = I_0 \text{sinc}^2\left(\frac{k a X}{2d}\right) \text{sinc}^2\left(\frac{k b Y}{2d}\right)$$

(b is length of slit in Y-dimension)

Plots of this in PPT file ~~are on the next~~
~~22 pages~~

(for $\frac{k}{2d} = 1$, just to see shape)

long + thin → \approx delta function in Y direction

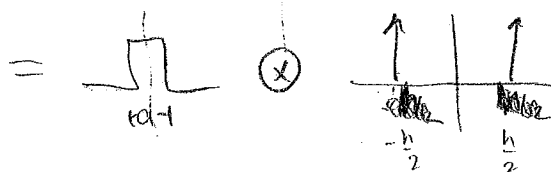
Application of the Convolution Theorem

The convolution theorem can make some problems trivial. For example:

Double Slit

1D Aperture Function

(ie other dimension = wide)



$$\mathcal{F}(arr) = \mathcal{F}(a) \times \mathcal{F}(rr)$$

$$= \frac{1}{\sqrt{2\pi}} a \operatorname{sinc} \frac{k_x a}{2} \times \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} (rr) e^{ik_x x} dx$$

as done previously
 $k_x = \frac{kx}{d}$

$$= \frac{1}{\sqrt{2\pi}} a \operatorname{sinc} \frac{k_x a}{2} \underbrace{\left(e^{-ik_x h/2} + e^{ik_x h/2} \right)}_{2 \cos \frac{k_x h}{2}}$$

$$\mathcal{F}(arr) = \frac{1}{\sqrt{2\pi}} 2a \operatorname{sinc} \frac{k_x a}{2} \cos \frac{k_x h}{2}$$

zeros/maxima of this probably given in 123

~~that's what I took from the book~~
~~that's what I remember the table!~~

$$I = I_0 \operatorname{sinc}^2 \frac{k_x a}{2} \cos^2 \frac{k_x h}{2}$$

double slit pattern!

Piece of cake!

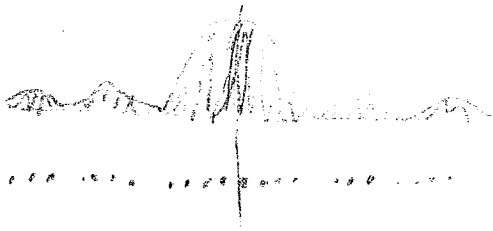
sinc^2
 \cos^2
(h > a so cos oscillates faster)

dm 35 pg 5

Discussion

Double Slit $I = I_0 \text{sinc}^2 \frac{kx a}{2d} \cos^2 \frac{kx b}{2d}$

Typically $a \ll b$



Maxima given by \cos^2 function

$$\frac{kx b}{2d} = m\pi$$

$$x = \frac{2d}{(2\pi/\lambda) b} m\pi = \frac{m d \lambda}{b}$$

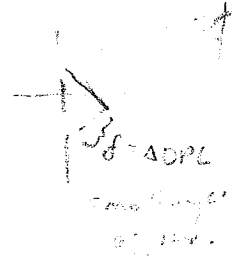
maxima $x = \frac{m d \lambda}{b}$

same as formula from 106 book

similarly

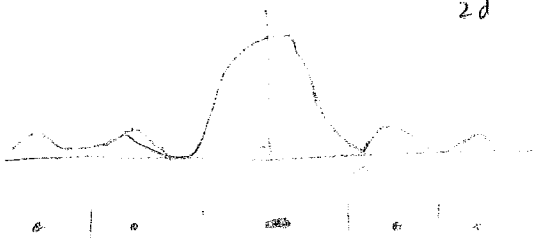
minima $x = \left(m + \frac{1}{2}\right) \frac{d \lambda}{b}$

Note: overall intensity modulated by sinc function



Single Slit

$$I = I_0 \text{sinc}^2 \frac{kx a}{2d}$$



minima same as sinc function (except $x=0$ maximum)

minima: $\frac{kx a}{2d} = m\pi$
 $x = \frac{2d}{(2\pi/\lambda) a} m\pi$

minima $x = \frac{m d \lambda}{a}$ $m \neq 0$

same as 106 book (hand-drawing argument there)

Note: overall intensity modulated by sinc function