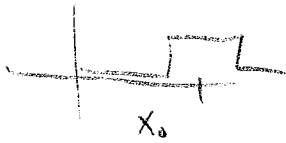


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FT of shifted function



$$\begin{aligned}
 FT(\text{---}\square L) &= FT(\text{---}\square L) \otimes FT(\text{---}\delta(x-x_0)) \\
 &= a \text{sinc} \frac{kx a}{2} \cdot \int_{-\infty}^{\infty} \delta(x-x_0) e^{ikx} dx \\
 &= a \text{sinc} \frac{kx a}{2} e^{ikx_0} \quad \text{(sifting property)}
 \end{aligned}$$

2D:  $FT(\text{shifted aperture}) = e^{ik_x x_0} e^{ik_y y_0} \cdot FT(\text{regular aperture})$

"Array Theorem"

$$FT(\text{array of identical apertures}) = FT(\text{aperture}) \cdot \sum_{\text{all apertures}} e^{i(k_x x_{0j} + k_y y_{0j})}$$

Ogan not multiply about factors of  $\chi(\text{off})!$

Don't ever have to be in an array!

$(x_{0j}, y_{0j})$ : center of  $j^{\text{th}}$  aperture  
 = FT of  $\begin{pmatrix} \uparrow & \uparrow & ? \\ \uparrow & \uparrow & \uparrow \text{ etc.} \end{pmatrix}$  (located at centers!)

Random vs Non-random array

↓

$\sum$  phase factors cancels

$I = I(\text{one aperture})$

$\sum$  phase factors gives you fine detail

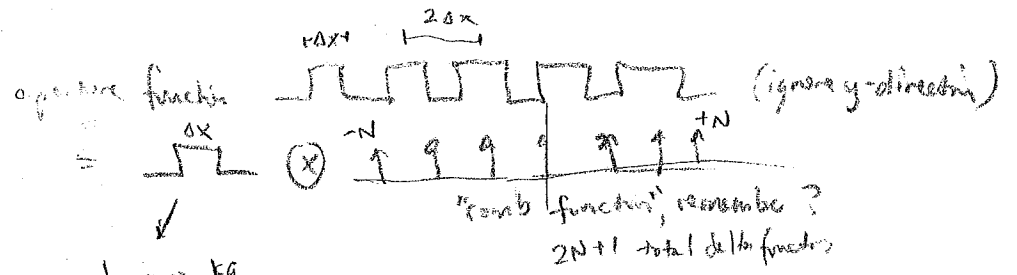
$I = I(\text{array})$  modulated by  $I(\text{one aperture})$

like  $\uparrow$  comes  $\times$   $\uparrow$  sinc<sup>2</sup>

for double slit example

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 2/11/17

Diffraction Gratings



Fourier Transform

$$\frac{1}{\sqrt{2\pi}} \text{sinc} \frac{kx \Delta x}{2}$$

$$+ \Delta x \text{sinc} \left( kx \frac{\Delta x}{2} \right)$$

$$FT = \frac{1}{\sqrt{2\pi}} \frac{\text{Si}((2N+1)\frac{\omega \Delta x}{2})}{\text{Si} \frac{\omega \Delta x}{2}}$$

from comb  
 $\omega \rightarrow 2kx$   
 $\omega \Delta x \rightarrow 2kx \Delta x$

let  $N$  now = total number

$$= \frac{1}{\sqrt{2\pi}} \frac{\text{Si} N k x \Delta x}{\text{Si} k x \Delta x}$$

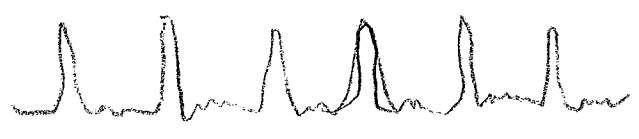
$$= \frac{1}{2\pi} \Delta x \text{sinc} \frac{kx \Delta x}{2} \frac{\text{Si} N k x \Delta x}{\text{Si} k x \Delta x}$$

$$I = I_0 \text{sinc}^2 \frac{kx \Delta x}{2} \frac{\text{Si}^2 N k x \Delta x}{\text{Si}^2 k x \Delta x}$$

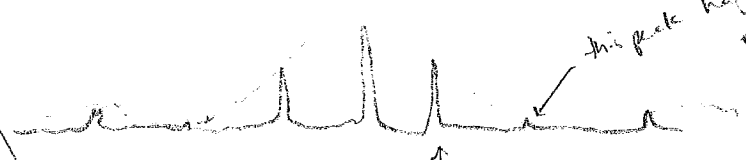
↑  
 added  $N^2$  in denominator (and removed  $\Delta x$  from Numer.)  
 So  $I(x=0) = I_0$

recall  $\frac{\text{Si} Nx}{\text{Si} x} \rightarrow N$  as  $x \rightarrow 0$   
 also from delta function bandwidth

Recall  $\frac{\text{Si} Nx}{x}$  plotted in delta function bandwidth for  $N=10$



this now is squared  
 & multiplied by  $\text{sinc}^2$



↑ this peak happens to fall on minima of  $\text{sinc}^2$  (because delta function spacing =  $2\Delta x$  exactly)

↑ position of this "first order peak" (and other peaks) depends on wavelength (via  $k$ )

"Fraunhofer" style  
 Multiple Slit  
 Approximation

of narrower aperture  
 wider  $\text{sinc}^2$

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section 11.8

Spectrometer

use (first order) diffraction spot to give you wavelength separation

let grating spacing =  $h$  now  
instead of  $2d \sin \theta$

- 1) position of peaks set by comb function
- 2) envelope of peaks set by sinc function

peaks occur when denominator goes to zero

$$k \times \frac{h/2}{d} = m\pi$$

$$\frac{2\pi}{\lambda} \times \frac{h/2}{d} = m\pi$$

$$\lambda = \frac{h}{m d}$$

$$\sin x = 0 \rightarrow x = m\pi$$

computer method  
 $m\lambda = d \sin \theta$

$x$  = position on screen

$d$  = distance from grating to screen

$h$  = grating spacing (microns/line)

Delta from delta function bandwidth

distance (in  $k$  space) to first zero =  $\frac{2\pi}{N \lambda_0}$

1/

$$\Delta k_x = \frac{2\pi}{N h}$$

$$\frac{x}{d} \Delta k = \frac{2\pi}{N h}$$

$$\frac{x}{d} \frac{2\pi d}{\lambda^2} = \frac{2\pi}{N h}$$

$$\Delta \lambda = \frac{d}{x} \frac{\lambda^2}{N h}$$

$$\frac{x h}{d} = m \lambda$$

$$= \frac{d}{m d} \frac{\lambda^2}{N}$$

$$\Delta \lambda = \frac{\lambda}{m N}$$

close to  $\Delta \lambda_{FWHM}$



$$RP = \frac{\lambda}{\Delta \lambda}$$

$$RP = m N$$

increase RP by increasing  $N$  (# lines illuminated by light) and/or  $m$  (but this decreases intensity)