# Optics: what you should already know, by Dr. Colton

#### From Phys 123 section 2

Complex numbers:

Euler's identity:  $e^{ix} = \cos x + i \sin x$ 

Complex numbers as points in the complex plane; polar  $\leftrightarrow$  rectangular conversion

General wave properties:

what all of these parameters mean: *x*, *t*, *A*,  $\lambda$ , *f*, *v*, *k*,  $\omega$ ,  $\phi$ 

 $f = A\cos(kx - \omega t + \phi) \leftrightarrow Ae^{i(kx - \omega t + \phi)}$ 

(and how to extend that to 3D for arbitrary wave direction and arbitrary oscillation direction)  $k = 2\pi/\lambda$ ;  $\omega = 2\pi/T$ 

 $v = \lambda f$ 

wave packets: 
$$v_{phase} = \omega/k$$
;  $v_{group} = \left(\frac{\partial \omega}{\partial k}\right)_{kave}$ 

Uncertainty relationships:

 $\Delta x \Delta k \geq \frac{1}{2} \; ; \; \; \Delta x \Delta p \geq \hbar/2$ 

$$\Delta t \Delta \omega \geq \frac{1}{2}$$
;  $\Delta t \Delta E \geq \hbar/2$ 

Reflection/transmission coefficients at normal incidence:

$$r = \frac{v_2 - v_1}{v_2 + v_1} = \frac{n_1 - n_2}{n_1 + n_2}; \quad t = \frac{2v_2}{v_2 + v_1} = \frac{2n_1}{n_1 + n_2}$$
$$R = |r|^2; \quad T = 1 - R$$

Fourier series:

$$f(x) = a_0 + \sum_{n=1}^{\infty} a_n \cos \frac{2\pi nx}{L} + \sum_{n=1}^{\infty} b_n \sin \frac{2\pi nx}{L}$$
$$(2\pi/L = k_0 = \text{fundamental [spatial] frequency})$$
$$a_0 = \frac{1}{L} \int_0^L f(x) dx$$
$$a_n = \frac{2}{L} \int_0^L f(x) \cos \frac{2\pi nx}{L} dx$$
$$b_n = \frac{2}{L} \int_0^L f(x) \sin \frac{2\pi nx}{L} dx$$

Fourier series in time:  $x \to t$ ;  $L \to T$ ;  $k_0 \to \omega_0$ 

Index of refraction, n

speed of light = c/n

 $\lambda_{\text{material}} = \lambda_{\text{vacuum}}/n$ Laws of reflection/refraction:

$$\theta_{incident} = \theta_{reflected}$$

$$\mathcal{F}_{incident} = \Theta_{reflected}$$

 $n_1 \sin \theta_1 = n_2 \sin \theta_2$  ( $\theta$  measured from the perpendicular)

Total internal refraction:  $\theta_{critical}$  of high index material is when  $\theta_2 = 90^{\circ}$  Polarization

Difference between linear and circular polarization

 $\theta_{\text{Brewster}} = \tan^{-1}(\theta_2/\theta_1)$ 

Difference between *s*- and *p*-polarization

## Lenses/mirrors

Thin lens equation: 1/f = 1/p + 1/q

Mirror: f = R/2

Lensmaker's eqn: 
$$\frac{1}{f} = (n-1)\left(\frac{1}{R_1} - \frac{1}{R_2}\right)$$

 $(R_1 = \text{pos}, R_2 = \text{neg if convex-convex})^2$ 

magnification: 
$$M = h_i/h_o = -q/p$$

f-number of a lens = f/D

Diffraction through slits/apertures:

Phase difference due to path-length difference:  $\phi = 2\pi (\Delta PL/\lambda)$ 

. .

Parallel ray approximation, if screen distance >> slit separation:  $\Delta PL = d\sin\theta$  (d = distance to reference of phase) Field at location on screen is sum of fields from each slit:  $E = E_0 (e^{i\phi 1} + e^{i\phi 2} + ...)$  (integrate if needed) Intensity  $I \sim |E|^2$ 

2 slit result:  $I = I_0 \cos^2\left(\frac{2\pi}{\lambda}\frac{d}{2}\sin\theta\right)$ ;  $d\sin\theta = m\lambda$  (maxima);  $d\sin\theta = (m + \frac{1}{2})\lambda$  (minima) 1 wide slit result:  $I = I_0 \operatorname{sinc}^2\left(\frac{\pi a \sin\theta}{\lambda}\right)$ ;  $a\sin\theta = m\lambda$  (minima) 2 wide slits result:  $I = (2 \text{ slit result}) \times (1 \text{ wide slit result})$ Arbitrary number/arrangement of slits: how to apply this technique to get  $I(\theta)$ Small angle approximation sometimes applies:  $\theta \approx \sin \theta \approx \tan \theta = y/L$ Circular aperture result, Rayleigh criterion:  $\theta_{\min, resolve} = 1.22 \lambda D$ Grating result:  $d\sin\theta_{\text{bright}} = m\lambda$ Spectrometer:  $R = \lambda_{ave}/\Delta\lambda = \#\text{slits} \times m$ 

Thin film interference:

 $OPL = PL \times n$  (PL = "path length"; OPL = "optical path length")

 $\triangle OPL + \text{other phase shifts} = m\lambda$  (constructive); ... =  $(m + \frac{1}{2})\lambda$  (destructive) Photons: (possibly not learned until Phys 222)

photon momentum  $p = h/\lambda$ 

photon energy  $E = pc = hc/\lambda$ 

### From Phys 220:

Coulomb's Law:

 $\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \, \hat{\mathbf{r}} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^3} \quad \text{(electric field from a point charge located at origin)}$  $\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q \, (\mathbf{r} - \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} \quad \text{(electric field from a point charge located at } \mathbf{r}')$ 

**Biot-Savart Law** 

 $\mathbf{B} = \frac{\mu_0}{4\pi} \int \frac{I \, d\ell' \times (\mathbf{r} - \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} \quad \text{(magnetic field from a current-carrying wire; integrate over the primed variables)}$ Gauss's Law (Maxwell #1):

 $\oint_{S} \mathbf{E} \cdot d\mathbf{a} = \frac{q_{enc}}{\varepsilon_{0}} \quad \text{(electric flux is proportional to q_{enclosed})}$ Gauss's Law for magnetism (Maxwell #2):

 $\oint_{S} \mathbf{B} \cdot d\mathbf{a} = \frac{q_{enc}}{\varepsilon_{0}} \quad \text{(no magnetic monopoles)}$ 

Faraday's Law (Maxwell #3):  $\oint \mathbf{E} \cdot d\boldsymbol{\ell} = -\frac{d\Phi_{\rm B}}{dt} \quad (\text{induced EMF is } -d(\text{flux})/\text{dt}; \text{ minus sign is Lenz's Law})$ Ampere's Law, with Maxwell correction (Maxwell #4):  $d\Phi_{\rm E} = -\frac{d\Phi_{\rm E}}{dt} = -\frac{d\Phi_{\rm E}}{dt} \quad (\text{induced EMF is } -d(\text{flux})/\text{dt}; \text{ minus sign is Lenz's Law})$ 

 $\oint \mathbf{B} \cdot d\boldsymbol{\ell} = \mu_0 I_{enc} + \epsilon_0 \mu_0 \frac{d\Phi_{\rm E}}{dt} \quad \text{(currents act as sources of magnetic fields; so do changing electric fields)}$ 

#### From Multivariable Calculus:

Scalar and vector functions:

f = a scalar function of x, y, z. Example:  $f(x, y, z) = x^2y + \sin z$ .  $\mathbf{A}$  = a vector function of x, y, z. Example:  $\mathbf{A}(x, y, z) = (x^2y + \sin z)\hat{\mathbf{x}} + xyz\hat{\mathbf{y}} + 4\hat{\mathbf{z}}$ , which means  $A_x = x^2y + \sin z$ ,  $A_y = xyz$ , and  $A_z = 4$ Gradient of a scalar function  $\nabla f = \frac{\partial f}{\partial x} \hat{\mathbf{x}} + \frac{\partial f}{\partial y} \hat{\mathbf{y}} + \frac{\partial f}{\partial z} \hat{\mathbf{z}}$  (which is a vector function)

Divergence of a vector function  $\nabla \cdot \mathbf{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$  (which is a scalar function) Curl of a vector function  $\nabla \times \mathbf{A} = \begin{vmatrix} \hat{\mathbf{x}} & \hat{\mathbf{y}} & \hat{\mathbf{z}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ A_x & A_y & A_z \end{vmatrix}$  (which is a vector function)

Laplacian of scalar function  $\nabla^2 f = \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2}$  (which is a scalar function)

What you should already know - pg 2

Laplacian of vector function  $\nabla^2 \mathbf{A} = \nabla^2 A_x \hat{\mathbf{x}} + \nabla^2 A_y \hat{\mathbf{y}} + \nabla^2 A_z \hat{\mathbf{z}}$  (which is a vector function) Handy "curl of curl" formula:  $\nabla \times (\nabla \times \mathbf{A}) = \nabla (\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}$ Gradient Theorem:

$$\int_{\substack{\text{path from}\\ \mathbf{a} \text{ to } \mathbf{b}}} (\nabla f) \cdot d\boldsymbol{\ell} = f(\mathbf{b}) - f(\mathbf{a})$$

Divergence Theorem:

$$\int_{volume} \nabla \cdot \mathbf{A} \, dv = \oint_{\substack{\text{surface bounding} \\ \text{the volume}}} \mathbf{A} \cdot d\mathbf{a}$$

Stokes' Theorem, aka Curl Theorem:

$$\int_{surface} (\nabla \times \mathbf{A}) \cdot d\mathbf{a} = \oint_{\substack{path \ bounding \\ the \ surface}} \mathbf{A} \cdot d\boldsymbol{\ell}$$