

July 32 (9)

Ch 10: Diffraction

Back to wave eqn: $\nabla^2 \vec{E} = \frac{1}{v^2} \frac{\partial^2 \vec{E}}{\partial t^2}$

harmonic time dependence, $\frac{d}{dt} \rightarrow -i\omega$
 $\nabla^2 \vec{E} = -k^2 \vec{E}$ (use $\frac{\omega}{v} = k$)
 "vector Helmholtz eqn"

rect: $\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$

spherical? $\vec{E} = \vec{E}(r, \theta, \phi)$

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \vec{E}}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \vec{E}}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \vec{E}}{\partial \phi^2}$$

simplest case: $E = E(r)$ only $\rightarrow = 0$

If not vector, $f = \frac{A}{r} \cos(kr - \omega t)$ works. I.e. $\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{df}{dr} \right) = \frac{1}{r} \frac{d^2 f}{dr^2}$
 $\nabla^2 E = -k^2 E$ "scalar Helmholtz eqn"

If vector that must satisfy Maxwell $\nabla \cdot \vec{E} = 0$ guess, since it's \perp to travel

$\vec{E} = \frac{A}{r} \cos(kr - \omega t) \hat{\phi}$ doesn't work! "At W P10.2 (assigned)

Can satisfy 1st three, but not divergence

$\vec{E}(r, \theta) = \frac{A \sin \theta}{r} \left(\cos(kr - \omega t) - \frac{1}{kr} \sin(kr - \omega t) \right) \hat{\phi}$ does work!

It's a pain to prove it ... in Griffith's III, but not P10.2

But... let's say $r = \text{big}$

$\theta \approx \pi/2$ true for travel close to z -direction since \vec{E} oscillation will then be in xy -plane

then that becomes $\vec{E} = \frac{A}{r} \cos(kr - \omega t) \hat{\phi}$

Discussion of scalar vs vector in "Significance of Scalar wave approx" on pg 261

Bottom line: For $r = \text{big}$ and $\theta \approx \pi/2$

then if E satisfies scalar Helmholtz, \vec{E} will satisfy (mostly) vector Helmholtz

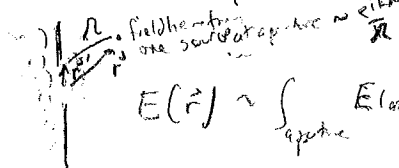
"scalar approximation"

Complex: $E = \frac{E_0 e^{ikr}}{r}$ or $E = E_0 \frac{e^{ikr}}{r}$ to give E_0 units of field

Where are we going?

Huygens: light as propagation of infinite number of spherical wavelets

Fig 10.1



$E(\vec{r}) \sim \int_{\text{aperture}} E(\text{at aperture}) \frac{e^{ikR}}{R} da'$

$\vec{R} = \vec{r} - \vec{r}'$
 $R = \sqrt{(x-x')^2 + (y-y')^2 + z^2}$

but: $n = R$

Can we derive this from the Wave Eqn? (or Helmholtz eqn)

Travel in z-direction

Expect $E = \frac{E_0}{r} e^{ikz}$

↑
let's make this more general,

Call it \tilde{E} , units of field

• expect something like a $\frac{1}{z}$ or $\frac{1}{r}$ dependence

Scalar Helmholtz:

$$\nabla^2 E = -k^2 E$$

Plug in guess

$$\nabla^2 (\tilde{E} e^{ikz}) = -k^2 (\tilde{E} e^{ikz})$$

Some derivatives...

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) (\tilde{E} e^{ikz}) = -k^2 \tilde{E} e^{ikz}$$

(product rule)

$$e^{ikz} \frac{\partial^2 \tilde{E}}{\partial x^2} + e^{ikz} \frac{\partial^2 \tilde{E}}{\partial y^2} + \frac{\partial}{\partial z} \left(\tilde{E} ik e^{ikz} + e^{ikz} \frac{\partial \tilde{E}}{\partial z} \right) = -k^2 \tilde{E} e^{ikz}$$

Another two product rules

$$ik \left[\tilde{E} ik e^{ikz} + e^{ikz} \frac{\partial \tilde{E}}{\partial z} \right] + \left[e^{ikz} \frac{\partial^2 \tilde{E}}{\partial z^2} + ik e^{ikz} \frac{\partial \tilde{E}}{\partial z} \right]$$

$$\left(-k^2 \tilde{E} + 2ik \frac{\partial \tilde{E}}{\partial z} + \frac{\partial^2 \tilde{E}}{\partial z^2} \right) e^{ikz}$$

All e^{ikz} 's cancel

$-k^2 \tilde{E}$ term cancels w/ RHS

Result: $\frac{\partial^2 \tilde{E}}{\partial x^2} + \frac{\partial^2 \tilde{E}}{\partial y^2} + 2ik \frac{\partial \tilde{E}}{\partial z} + \frac{\partial^2 \tilde{E}}{\partial z^2} = 0$

$\Delta \approx 0$

"paraxial wave approximation"
Amplitude of wave varies slowly
in z direction so it
looks much like a plane wave

The or maybe "a" Soln to this eqn:

HW P10.5 (unassigned)

$$\tilde{E} = -\frac{i}{\lambda z} \iint E(x', y') e^{i \frac{k}{2z} [(x-x')^2 + (y-y')^2]} dx' dy'$$

do integral, then $E = \tilde{E} e^{ikz}$
 aperture where $z=0$
 field strength across aperture (if varying) ← "aperture function"

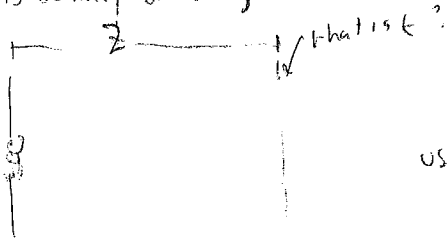
gives answer just like original guess with

$$R = [(x-x')^2 + (y-y')^2 + z^2]^{1/2} \\ = z \left[1 + \frac{(x-x')^2}{z^2} + \frac{(y-y')^2}{z^2} \right]^{1/2} \\ \approx z \left(1 + \frac{1}{2} \frac{(x-x')^2}{z^2} + \frac{1}{2} \frac{(y-y')^2}{z^2} + \dots \right)$$

but with proper $-\frac{i}{\lambda}$ out in front

just keep zeroth order for denom. ($\frac{1}{z}$)
 • keep first order for exponent (need that since small change in $x \rightarrow$ big change in e^{ikx})

To actually do integral



use $(x-x')^2 = x^2 - 2xx' + x'^2$
 ↓
 can be taken out of integral

(same for y)

$$\star E = -\frac{i}{\lambda z} e^{ikz} e^{i \frac{k}{2z} (x^2 + y^2)} \iint_{\text{aperture}} E(x', y') e^{i \frac{k}{2z} (x'^2 + y'^2)} dx' dy'$$

"The Fresnel Approximation" Eqn. 10.13

Babinet: if aperture then $\iint_C f(x) dx = \iint_{\text{aperture}} f(x) dx$
 (square lets light through gets interference)

Not so interesting?

diffraction from hair $\frac{1}{\text{hair}} = \iint_{\text{aperture}} - \iint_{\text{slit}}$ → hair diffraction very similar to slit!

from → diffraction like = (wool's salad)
 Heile's circle bright spot inevitable! "Poisson's spot"

Note: There are 4 complicated integral formulas in this chapter.

I only care about 2 of them really (except problem P10.6)

* "Fresnel Approx." Eqn 10.13, that we just discussed

* "Fresnel's Diffraction Formula" Eqn 10.1

$$E(x, y, z) = \frac{-i}{\lambda} \iint_{\text{aperture}} E(x', y', 0) \frac{e^{ik\sqrt{(x-x')^2 + (y-y')^2 + z^2}}}{\sqrt{(x-x')^2 + (y-y')^2 + z^2}} dx' dy'$$

Developed before Maxwell's eqns. Not correct.
"Gessed"

* "Fresnel-Kirchhoff Diffraction formula" Eqn 10.10

Same thing, w/ additional term

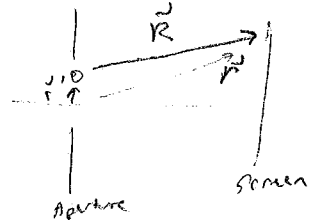
in integrand: $\frac{1 + \cos \theta}{2}$

$\theta = \angle$ between \vec{R} and z -axis

$$\vec{R} = \vec{F} - \vec{F}'$$

- Correct, but hard to use.

- Turns into Fresnel Approx, Eqn 10.13 in approximations we made earlier.



* "Fraunhofer Approx" Eqn 10.19

which we'll talk about soon.

Notes on the P10.6

→ see PPT slide

Babinet's

If aperture =



light through space,
blocked by circle

$$\text{then } \iint_{\square} \dots da' = \iint_0 \dots da'$$

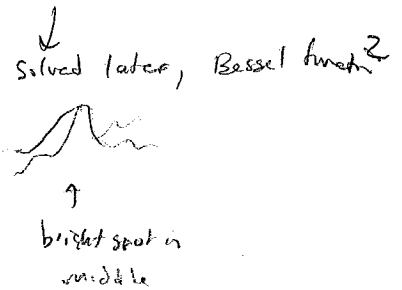
Not so interesting?

Diffraction from hair:

$$E = \iint_{\text{hair}} = \underbrace{\iint_{\text{all space}}}_{\text{no features}} - \iint_{\text{slit}}$$

$$|E_{\text{from hair}}|^2 \text{ similar to } |E_{\text{single slit}}|^2$$

I from blocked circle like I open circle



↓
bright spot in middle?

Judge: Simon Poisson (1818)
same as Poisson distribution

"No, can't be!"

Francis Arago measured

"Yes!"

Fresnel wins award!