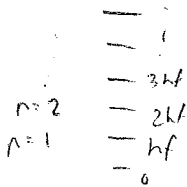


Blackbody

"Hot objects glow"

Boltzmann factor: Prob of an energy state $\propto e^{-E/k_B T}$

Planck: light out of particles frequency f , only available in integral multiples of hf



Today: $hf = \text{photon energy}$

Planck: "act of desperation" to fit experimental data

(Skipping the UV catastrophe and Rayleigh-Jeans incorrect formula. Just going to correct description.)

Prob of being in n^{th} state = $\frac{e^{-E_n/k_B T}}{\sum \text{all B.F.s}}$ (no chemical potential μ)

$$P_n = \frac{e^{-n hf/k_B T}}{e^0 + e^{-hf/k_B T} + e^{-2hf/k_B T} + \dots}$$

$$= e^{-n hf/k_B T} \underbrace{\left(1 - e^{-hf/k_B T}\right)}_{\text{result of infinite sum}}$$

actual energy in n^{th} state = $(n hf) \times P_n$

Total energy of all f -related states

$$E = \sum_{n=0}^{\infty} (n hf) \left[e^{-n hf/k_B T} \left(1 - e^{-hf/k_B T}\right) \right]$$

$$= hf \left(1 - e^{-hf/k_B T}\right) \sum_{n=0}^{\infty} n e^{-n hf/k_B T}$$

Trick: $\sum_{x=0}^{\infty} x e^{-ax} = - \sum_{x=0}^{\infty} \frac{\partial}{\partial a} (e^{-ax})$

$$= - \frac{\partial}{\partial a} \sum_{x=0}^{\infty} e^{-ax} = - \frac{\partial}{\partial a} \left(\frac{1}{1 - e^{-a}} \right)$$

* or sum with Mathematica

$$= hf \left(1 - e^{-hf/k_B T}\right) \left[\frac{e^{-hf/k_B T}}{\left(1 - e^{-hf/k_B T}\right)^2} \right]$$

$$\frac{e^{-a}}{(1 - e^{-a})(1 - e^{-a})} = \frac{e^{-a}}{(1 - e^{-a})^2}$$

$$= \frac{hf}{e^{hf/k_B T} - 1}$$

Cancel, with (-)

$$\frac{1}{(e^{hf/k_B T} - 1)(1 - e^{-hf/k_B T})}$$

$$E = \frac{hf}{e^{hf/k_B T} - 1}$$

energy associated with all photons frequency f or plane waves with that frequency

"Box" (could be $L \rightarrow \infty$)



radiation inside \rightarrow which frags are possible?

We'll find
not all frags
are equally likely

Body rad



Integral # of half-wavelengths, $L_x = n \frac{\lambda_x}{2} \rightarrow \lambda_x = \frac{2L}{n}$

Allowed k 's: $k_x = \frac{2\pi}{\lambda_x}$

$k_x = \frac{\pi}{L} \times n = k_0 n$

where $k_0 = \frac{\pi}{L}$

fact of 2 diff. from P & W

similarly

$k_y = k_0 m$

$k_z = k_0 l$

since this is only $e^{i(kx - \omega t)}$ form expansion, not standing waves

$k = \sqrt{k_x^2 + k_y^2 + k_z^2} = k_0 \sqrt{n^2 + m^2 + l^2}$

$= k_0 S$

$S = \sqrt{n^2 + m^2 + l^2}$

(like $|\vec{r}|$ for 3D space)

$\frac{\omega}{k} = c \rightarrow k \text{ also} = \frac{\omega}{c} = \frac{2\pi f}{c}$

so $f = \frac{c}{2\pi} k_0 S$

$U_{rad} = 2 \sum_{\substack{\text{modes} \\ \text{polar.}}} \frac{hf}{e^{hf/kT} - 1}$

\leftarrow Assumes each allowed $k_x, k_y, k_z =$ equally probable

$= 2 \sum_{\substack{\text{mode} \\ = 0}}^{\infty} \frac{h \left(\frac{c}{2\pi} k_0 S \right)}{e^{\frac{h \left(\frac{c}{2\pi} k_0 S \right)}{kT}} - 1}$

\swarrow weighted towards higher S
 \rightarrow higher (number)
 \rightarrow higher (frequency)
 \rightarrow More of the allowed modes can give you higher freq. than lower freq.

$= 2 \frac{hc}{2\pi} k_0 \frac{1}{8} \int_0^{\infty} \frac{S}{e^{\frac{hc k_0 S}{2\pi kT}} - 1} \times 4\pi S^2 dS$

\downarrow only positive n, m, l
 $(\frac{1}{2})^3$ diff. from P & W

$x = \frac{hc k_0 S}{2\pi kT}$

$S = \frac{2\pi kT}{hc k_0} x$
 $dS = \left(\frac{2\pi kT}{hc k_0} \right) dx$

$2 \frac{hc}{2\pi} k_0 \frac{4\pi}{8} \left(\frac{2\pi kT}{hc k_0} \right)^4 \int_0^{\infty} \frac{x^3 dx}{e^x - 1}$

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$$\frac{1}{2} \frac{K c^3}{h^3 c^3} \frac{16 \pi^4 k^4 T^4}{h^3 c^3 k_0} \frac{\pi^4}{15}$$

$$\frac{8}{15} \frac{\pi^8}{h^3 c^3} \frac{k^4 T^4}{\pi^3 / L^3}$$

$$\frac{8}{15} \frac{\pi^5 L^3 k^4 T^4}{h^3 c^3}$$

$$\pi = \frac{L}{2r}$$

$$h = 2r$$

$$u = \frac{U}{L^3} = \frac{8}{15} \frac{\pi^5 k^4}{h^3 c^3} T^4$$

Planck Law 13.2

Section 13.1

$$I = \frac{c}{4} \cdot u$$

I = total intensity of radiation flowing outwards

c from speed of light moving away

$\frac{1}{2}$ = half energy travelling away, half towards

$\frac{1}{4}$ energy from $\frac{1}{2}$ is distributed over $\frac{1}{2}$ area

with "intensity observed along a path will be actual intensity times cosine"

So

$$I = \frac{2}{15} \frac{\pi^5 k^4}{h^3 c^2} T^4$$

$$\sigma = 5.67 \cdot 10^{-8} \frac{W}{m^2 K^4}$$

Stefan Boltzmann constant

collecting factor emissivity (from both sides surface radiates)

$$I = \epsilon \sigma T^4 \quad \text{Stefan-B Law}$$

How much of u is at each frequency?

ρ_f

energy density in f

$$\text{ie } u = \int \rho_f df$$

Go back to integral

$$u = \frac{2}{2\pi} \frac{hc}{k_0} \frac{1}{8} \int_0^\infty \frac{5}{e^{\frac{hc k_0 s}{2\pi k T}} - 1} 4\pi s^2 ds$$

Go back to frequency: substitute in $f = \frac{s}{2\pi} k_0 s$

$$\text{Simplify } u = \int_0^\infty \rho_f df$$

$$\text{with } \rho_f = \frac{8\pi h f^3}{c^3 (e^{hf/kT} - 1)}$$

"Planck radiation law"

ρ_f



blackbody spectrum