

Lydane-Sachs-Teller

back to Sellmeier, but w/ $\gamma=0$

$$\epsilon = \epsilon_{\infty} + \frac{\epsilon_0 - \epsilon_{\infty}}{1 - \frac{\omega^2}{\omega_T^2}}$$

when $\epsilon=0$, $\omega = \omega_L$

$$0 = \epsilon_{\infty} + \frac{\epsilon_0 - \epsilon_{\infty}}{1 - \frac{\omega_L^2}{\omega_T^2}}$$

\uparrow
 x

$$0 = A + \frac{B-A}{1-x}$$

$$0 = A - Ax + B - A$$

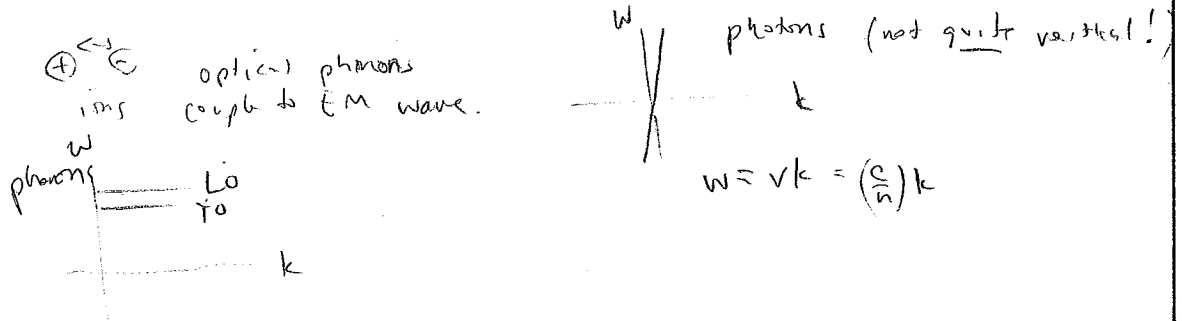
$$x = \frac{B}{A}$$

$$\left(\frac{\omega_L}{\omega_T}\right)^2 = \frac{\epsilon_0}{\epsilon_{\infty}}$$

LST

Polaritons - coupling of em wave w/ something

Phonon polaritons (the only type mentioned in Kittel)

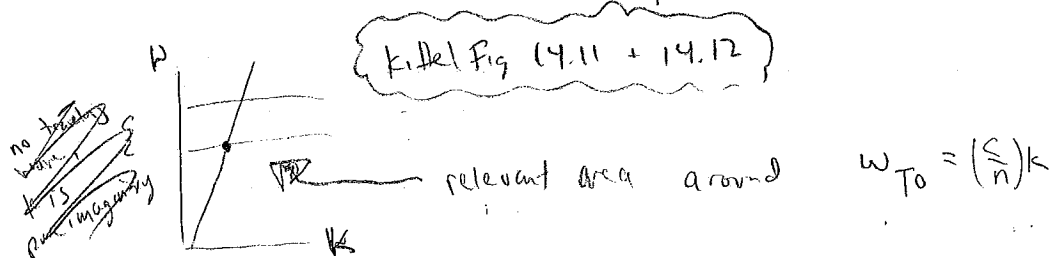


In regions where frequencies are degenerate (ω_0),
really need degenerate pert. theory (like we did
for band structure ~~to~~ when two electron
modes had same frequency).

Result \rightarrow two coupled modes, "upper branch"
+ "lower branch"

At frequencies farther away, looks regular again.

Instead of QM degenerate pert. theory, use this method which gives
essentially correct result.



Force dielectric constants to be the same

$$\text{Light: } \frac{\omega}{k} = \frac{c}{n} = \frac{c}{\sqrt{\epsilon}} \rightarrow \epsilon = \left(\frac{ck}{\omega}\right)^2$$

$$\text{TO phonons: } \epsilon = \epsilon_{\infty} + \frac{\epsilon_0 - \epsilon_{\infty}}{1 - \frac{\omega^2}{\omega_T^2}} \quad (\text{no damping for simplicity})$$

$$\frac{c^2 k^2}{\omega^2} = \epsilon_{\infty} + \frac{\epsilon_0 - \epsilon_{\infty}}{1 - \frac{\omega^2}{\omega_T^2}}$$

polariton dispersion relation

can in theory solve ω vs k

$\times \omega^2$

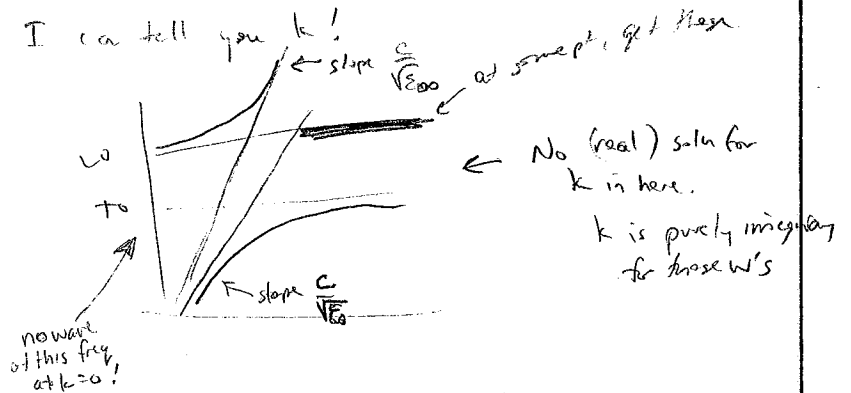
$$c^2 k^2 = \epsilon_{\infty} \omega^2 + \frac{\omega^2 (\epsilon_0 - \epsilon_{\infty})}{1 + \frac{\omega^2}{\omega_T^2}}$$

Hard to solve for ω ,

Easier for k

$$k = \frac{1}{c} \sqrt{\epsilon_{\infty} \omega^2 + \frac{\omega^2 (\epsilon_0 - \epsilon_{\infty})}{1 + \frac{\omega^2}{\omega_T^2}}}$$

For a given ω , I can tell you k !



Consider the lower branch, small k / small ω

$$k = \frac{1}{c} \sqrt{\epsilon_{\infty} \omega^2 + \frac{\omega^2 (\epsilon_0 - \epsilon_{\infty})}{1 + \text{small}}}$$

$$k = \frac{1}{c} \omega \sqrt{\epsilon_0}$$

Consider upper branch, large k/ω

$$k = \frac{1}{c} \sqrt{\epsilon_{\infty} \omega^2 + \frac{\omega^2 (\epsilon_0 - \epsilon_{\infty})}{\cancel{\omega^2} + \frac{\omega^2}{\omega_T^2}}}$$

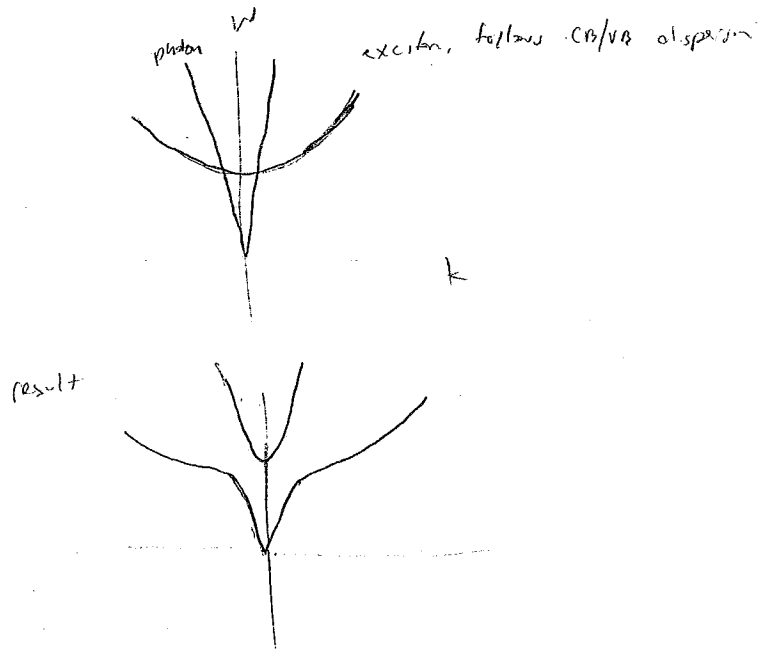
$$\sqrt{\epsilon_{\infty} \omega^2 + \frac{\omega^2 (\epsilon_0 - \epsilon_{\infty})}{\frac{\omega^2}{\omega_T^2}}}$$

then this small

$$k = \frac{1}{c} \omega \sqrt{\epsilon_{\infty}}$$

Exciton polariton

Do excitons interact w/ EM wave? of course!



Plasmons
- back to
- conductivity

$$\epsilon = 1 + \frac{ne^2}{m\epsilon_0 \omega_0^2 - \omega^2}$$

\uparrow ω_p^2 \uparrow
 No restoring force now!
 No resonant freq!

$$\epsilon = 1 - \frac{\omega_p^2}{\omega^2}$$

or really

$$\epsilon = \epsilon_{\infty} - \frac{\omega_p^2}{\omega^2}$$

$$= \epsilon_{\infty} \left(1 - \frac{\omega_p^2 / \epsilon_{\infty}}{\omega^2} \right)$$

$$= \epsilon_{\infty} \left(1 - \frac{\omega_{p,eff}^2}{\omega^2} \right)$$

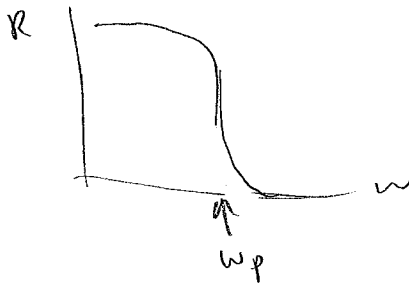
if "positive" core background has dielectric constant ϵ_{∞} essentially constant to frequencies well above ω_p

Like w/ LO phonon, this has a longitudinal resonance when $\epsilon = 0$

When $\omega < \omega_p$
 $\epsilon < 0$

$n = \sqrt{\epsilon} = \text{purely imaginary!}$

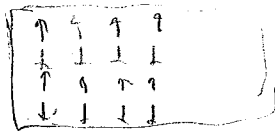
then $R = \left| \frac{1 - \tilde{n}}{1 + \tilde{n}} \right|^2 = 100\% \text{ (can show)}$



Now mention

At ω_p , $\epsilon = 0$, \rightarrow longitudinal resonance.

Kittel Fig 14.4



arrows = displacement of electrons

"plasmon" - a quantized plasma oscillation
com in units of $\hbar \omega_p$