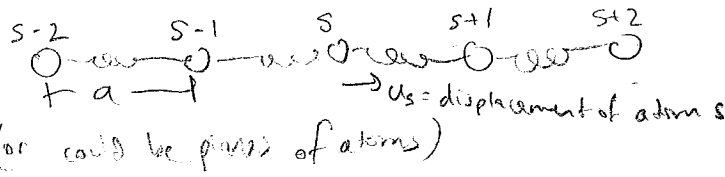


Phase vs Group Velocity

Fresh model: For now all identical items



"Harmonic Approximation"

Newton 2:  $\sum F_s = m a$

$$\underbrace{C(u_{s+1} - u_s)}_{\substack{\text{spring} \\ \text{constant} \\ \text{force from right-hand spring}}} - \underbrace{C(u_s - u_{s-1})}_{\text{left hand}} = m \frac{d^2 u_s}{dt^2}$$

$$m \frac{d^2 u_s}{dt^2} = C [u_{s+1} - 2u_s + u_{s-1}]$$

Guess  $u \approx e^{i(kx - \omega t)}$

Going from  $u_s$  to  $u_{s+1} \rightarrow \Delta x = a \rightarrow$  phase factor of  $e^{ika}$

$u_{s+1} = e^{ika} u_s$

Plug in to eqn:

$$m(-\omega^2) u_s = C [e^{ika} - 2 + e^{-ika}] u_s$$

$e^{ika} + e^{-ika} = 2 \cos ka$

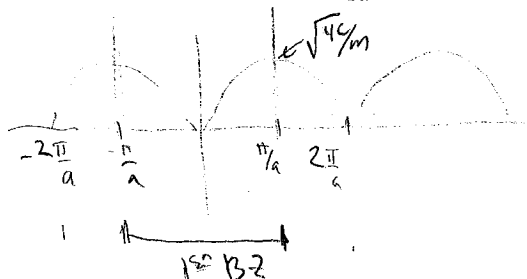
$$m(\omega^2) = 2C [1 - \cos ka]$$

$$\boxed{\omega = \left[ \frac{2C}{m} (1 - \cos ka) \right]^{1/2}}$$

dispersion relation

$$\boxed{\omega = \sqrt{\frac{4C}{m}} \left| \sin \frac{ka}{2} \right|}$$

because  $\sin^2 \frac{x}{2} = \frac{1}{2} (1 - \cos x)$



initial slope?

$$\frac{d\omega}{dk} = \sqrt{\frac{4c}{m}} \cos \frac{ka}{2} \left(\frac{a}{2}\right) \Big|_{\text{small } k}$$

$$v = \frac{a}{2} \sqrt{\frac{4c}{m}} = a \sqrt{\frac{c}{m}}$$

informative? starts out linear, not linear.

Periodicity? repeats over  $\frac{ka}{2} = n\pi$  ( $\sin x$  has period  $\pi$ , not  $2\pi$ )

$$k \approx \frac{2\pi}{a} n$$

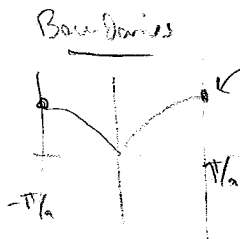
So  $k' = k + \frac{2\pi n}{a}$  will have same value as  $k$

What's special about

$\frac{2\pi}{a}, \frac{4\pi}{a}, \frac{6\pi}{a}$  etc?

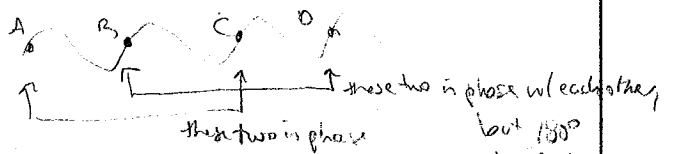
Reciprocal Lattice Vectors!

Summary: by subtracting an appropriate RLV from  $k$ , we always obtain an equivalent wavevector in 1st BZ



flat slope  $\rightarrow v = 0$   
Does that make sense?

Yes:  $k = \pm \frac{\pi}{a} \rightarrow \frac{2\pi}{\lambda} = \frac{\pi}{a} \rightarrow \lambda = 2a$



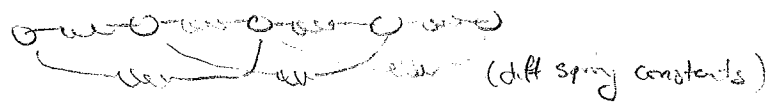
↑ ↓ ↑ ↓

↑ ↓  
↓ ↑

This is a standing wave! Of course no velocity.

day 13 pg 3

Possible extension forces between next nearest neighbors



or between 3<sup>rd</sup> nearest neighbors -- or 4<sup>th</sup>, etc.

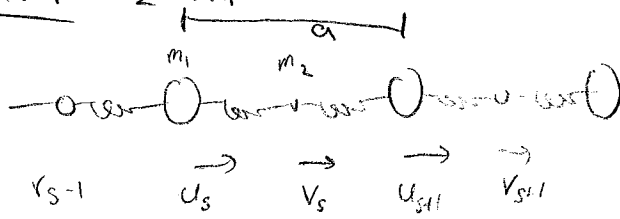
result:  $C_p$  = spring constant between  $p^{\text{th}}$  nearest neighbors

$$\omega = \left[ \frac{2}{m} \sum_{p=1}^{\infty} C_p (1 - \cos pka) \right]^{1/2}$$

back + shows inverse relationship is

$$C_p = -\frac{M a}{2\pi} \int_{-\pi/a}^{\pi/a} \omega^2 \cos pka \, dk$$

Next model: 2 different atoms



$$m_1 \frac{d^2 u_s}{dt^2} = c \left[ (v_s - u_s) - (u_s - v_{s-1}) \right]$$

$\downarrow -\omega^2 u_s$ 
 $\downarrow v_s e^{-ika}$

$$+\omega^2 m_1 u_s = +2c u_s + c(1 + e^{-ika}) v_s$$

Similarly  $+ \omega^2 m_2 v_s = -c(1 + e^{+ika}) u_s + 2c v_s$  coupled equations!

Solve w/ Matrix:  $0 = \begin{pmatrix} 2c - \omega^2 m_1 & -c(1 + e^{-ika}) \\ -c(1 + e^{+ika}) & 2c - \omega^2 m_2 \end{pmatrix} \begin{pmatrix} u_s \\ v_s \end{pmatrix}$

type is Eqn 21 pg 98!

$\hookrightarrow \det = 0$  for non-trivial soln

$$(2c - \omega^2 m_1)(2c - \omega^2 m_2) - c^2 (1 + e^{+ika})(1 + e^{-ika}) = 0$$

$$1 + e^{ika} + e^{-ika} + 1 = 2 + 2\cos ka$$

$$= 2c^2(1 + \cos ka)$$

$$m_1 m_2 \omega^4 + 2c(m_1 + m_2)\omega^2 + c^2(4 - 2(1 + \cos ka)) = 0$$

$$4 - 2 - 2\cos ka$$

$$2c^2(1 - \cos ka) = 0$$

quadratic eqn in  $\omega^2$

$$\omega^2 = \frac{+2c(m_1 + m_2) \pm \sqrt{4c^2(m_1 + m_2)^2 - 4(m_1 m_2)(2c^2(1 - \cos ka))}}{2m_1 m_2}$$

$\omega = \sqrt{\dots}$   $\rightarrow$  All that

plot w/ Mathematica for

$m_1 = 3, m_2 = 1, c = 1, a = 1$

Hand out to class

# Linear chain of balls + springs, 2 different atoms

In[16]:= m1 = 3; m2 = 1; c = 1; a = 1;

In[17]:= w1[k\_] = Sqrt[2 c (m1 + m2) + Sqrt[4 c^2 (m1 + m2)^2 - 4 (m1 m2) (2 c^2 (1 - Cos[k a]))]]]

Out[17]=  $\sqrt{8 + \sqrt{64 - 24 (1 - \cos[k])}}$

In[18]:= w2[k\_] = Sqrt[2 c (m1 + m2) - Sqrt[4 c^2 (m1 + m2)^2 - 4 (m1 m2) (2 c^2 (1 - Cos[k a]))]]]

Out[18]=  $\sqrt{8 - \sqrt{64 - 24 (1 - \cos[k])}}$

In[19]:= Plot[{w1[k], w2[k]}, {k, -Pi, Pi}]

