

day 22 pg 1

Some numbers...

$a = 4.225 \text{ \AA}$

Sodium
1 valence electron
atom

$$n = \frac{N_{\text{electrons}}}{V} = \frac{N_{\text{atoms}}}{V} = 2.652 \cdot 10^{28} / \text{m}^3 \text{ from Table 1.4, (b) 21}$$

Let's say $m_{\text{electron}} = 9.11 \cdot 10^{-31} \text{ kg}$ (actually we'll find that m_{eff} acts like it's reduced. It doesn't require as much force to produce acceleration as you might think)

$$\text{Then } E_F = \left(3n \right)^{2/3} \cdot \left(2.652 \cdot 10^{28} \text{ m}^{-3} \right)^{2/3} \cdot \frac{\left(1.055 \cdot 10^{-34} \text{ Js} \right)^2}{2 \cdot 9.11 \cdot 10^{-31} \text{ kg}}$$

$$= 5.199 \cdot 10^{-19} \text{ J} \times \frac{1 \text{ eV}}{1.6 \cdot 10^{-19} \text{ J}}$$

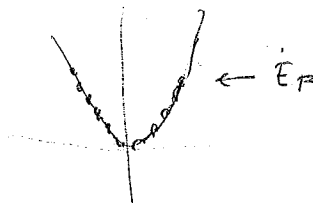
$$= \boxed{3.25 \text{ eV}}$$

compare to $k_B T = (1.38 \cdot 10^{-23}) (300) \text{ at room T}$

$$= 4.14 \cdot 10^{-21} \text{ J}$$

$$= \boxed{0.0258 \text{ eV}}$$

the assumption of T close to 0 is justified!



only electrons very close to here can even absorb kT of energy (deep electrons can't move up in energy because nearby states are filled)

$$E^2 = \frac{p^2 \hbar^2}{2m} = \frac{\hbar^2 k^2}{2m}$$

$$k_F = \frac{1}{\hbar} \sqrt{2m E_F}$$

$$= \frac{1}{1.055 \cdot 10^{-34}} \sqrt{2 \cdot 9.11 \cdot 10^{-31} \cdot 5.199 \cdot 10^{-19}}$$

$$= 9.23 \cdot 10^9 \text{ m}^{-1} \approx \frac{1 \text{ m}}{10^{10} \text{ \AA}}$$

$$\boxed{k_F = 0.923 \text{ \AA}^{-1}}$$

compare β edge

$$k_{\text{edge}} \frac{\pi}{a} = \frac{\pi}{4.225 \text{ \AA}} = \boxed{0.7448 \text{ \AA}^{-1}}$$

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$$\frac{V}{2\pi^2} \left(\frac{2m}{\hbar^2} \right)^{3/2} \left[\frac{\hbar^2}{2m} \left(3\pi^2 N \right)^{2/3} \right]^{1/2} E_F^{1/2}$$

$$\frac{V}{2\pi^2} \left(\frac{2m}{\hbar^2} \right)^{3/2} \frac{\hbar}{\sqrt{2m}} \left(3\pi^2 N \right)^{1/3} E_F^{1/2}$$

$$\frac{V}{2\pi^2} \left(\frac{2m}{\hbar^2} \right)^{3/2} E_F^{1/2}$$

Also $E_F = \frac{\hbar^2}{2m} \left(3\pi^2 N \right)^{2/3}$

$D(E_F) = \frac{V}{2\pi^2} \left(\frac{2m}{\hbar^2} \right)^{3/2} E_F^{1/2} \rightarrow$ want to simplify

Also $E_F = \frac{\hbar^2}{2m} \left(3\pi^2 N \right)^{2/3}$ from yesterday or earlier today

↓ solve for V

$$\left(\frac{2m E_F}{\hbar^2} \right)^{3/2} = 3\pi^2 N \frac{V}{V}$$

$$V = \left(3\pi^2 N \left(\frac{2m E_F}{\hbar^2} \right)^{-3/2} \right)$$

"and then a mistake occurs"

$$D(E_F) = \left[3\pi^2 N \left(\frac{2m E_F}{\hbar^2} \right)^{-3/2} \right] \left(\frac{2m}{\hbar^2} \right)^{3/2} E_F^{1/2}$$

$$D(E_F) = \frac{3}{2} \frac{N}{E_F}$$

useful simplified form as you'll see in a minute

But this is not a function of E that is $D(E) \stackrel{not}{\sim} \frac{1}{E}$

end of step (c)

Back to calculation of U

$$U = \int E D(E) f(E) dE$$

approximate this as $D(E_F)$, take out of integral

$\frac{1}{e^{+(E-E_F)/kT} + 1}$

approximate this as E_F , a constant

$$U \approx D(E_F) \int_0^{\infty} \frac{E dE}{e^{(E-E_F)/kT} + 1}$$

end of step (d)

this works because f changes most rapidly around E_F , so when we take $\frac{\partial U}{\partial T}$ later, the only non-zero part will be when $E \approx E_F$

$C = \frac{\partial U}{\partial T} \rightarrow$ take $\frac{\partial}{\partial T}$ inside integral

$$\frac{\partial}{\partial T} \left(\frac{1}{e^{(E-E_F)/kT} + 1} \right) = \frac{1}{(e^{(E-E_F)/kT} + 1)^2} \times e^{(E-E_F)/kT} \cdot \left(-\frac{(E-E_F)}{kT^2} \right)$$

$$C = D(E_F) \frac{1}{kT^2} \int_0^{\infty} \frac{(E-E_F) E e^{(E-E_F)/kT}}{(e^{(E-E_F)/kT} + 1)^2} dE$$

let $x = \frac{E-E_F}{kT} \rightarrow E = kTx + E_F$
 $dE = kT dx$

$$C = D(E_F) k \int_{-E_F/kT}^{\infty} \frac{x e^x}{(e^x + 1)^2} (kTx + E_F) dx$$

end of step (e)

small T approx: take integral from $-\infty$ to $+\infty$, then this part integrates to 0 since $\frac{x e^x}{(e^x + 1)^2}$ is odd

$$C = D(E_F) k T \int_{-\infty}^{\infty} \frac{x^2 e^x}{(e^x + 1)^2} dx$$

$\hookrightarrow \pi^2/3$

$$= \left(\frac{3}{2} \frac{N}{E_F} \right) (kT) \left(\frac{\pi^2}{3} \right)$$

$$C_{el} = \frac{\pi^2}{2} \frac{k^2 T N}{E_F}$$

free electrons

$C \sim T$

Combined w/ phonons

$$C = \underbrace{\gamma T}_{\text{electron}} + \underbrace{AT^3}_{\text{phonon}} \quad \text{at low } T$$

$$\frac{C}{T} = \gamma + AT^2 \quad \left. \begin{array}{l} \text{y axis} \uparrow \\ \text{x axis} \end{array} \right\} \text{Fig 9 for potassium pg 145}$$

$$\gamma = b + mx \quad \text{fits perfectly!}$$

Table 2 pg 146 - Experimental γ vs Theoretical γ

→ pretty close, but some differences.

Attributed difference to " (thermal) effective mass"

Trace back --- where's the m?

$$C \sim D(E_F)$$

$$C \sim m^{3/2} E_F^{1/2}$$

$$\hookrightarrow E_F \sim \frac{1}{m}, \text{ so this } \sim \frac{1}{m^{1/2}}$$

$$C \sim m$$

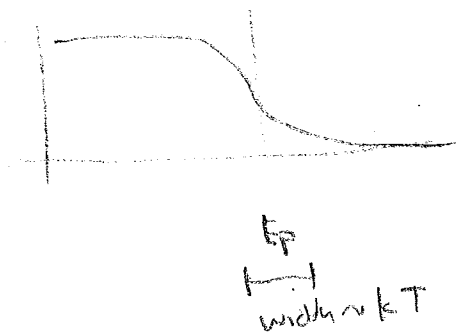
Check on this

maybe $C \sim \frac{1}{m}$

$$\text{Since } E_F = \frac{\hbar^2 k^2}{2m} \left(\frac{3N}{4\pi} \frac{2\pi}{v} \right)^{2/3}$$

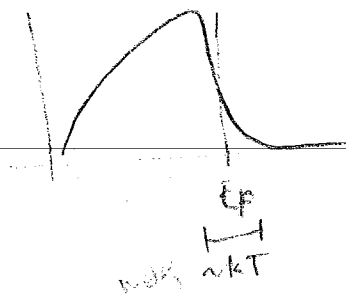
Qualitative calculation of C_v

At low T 's 1) $f(E)$ looks like



2) $D(E) \propto \sqrt{E}$

Multiply together



how many states there? (that are involved in "action")

$$= N \times \frac{kT}{E_F}$$

if each one gains kT ,

then

$$U = N \cdot \frac{k^2 T^2}{E_F}$$

$$C_v = \frac{2 N k^2 T}{E_F}$$

"Fermi surface"

→ more in Ch 9

(which we will not really discuss)

compare to $\frac{1}{2} N \frac{k^2 T}{E_F}$
 constant at all T 's
 right on!

out of town Mon & Wed