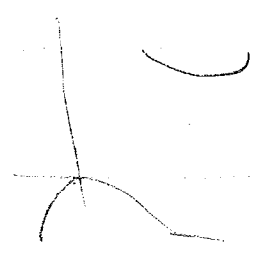


Ch 8  
Effektive Mass



real band gap opens up due to crystal potential  
- near zone edges, it's quadratic.

← put electron at bottom of slope, how does it behave?

Free electron  $E = \frac{\hbar^2 k^2}{2m} \rightarrow \frac{dE}{dk} = \frac{\hbar^2}{m} k$   
 $\frac{d^2E}{dk^2} = \frac{\hbar^2}{m}$

Electron in crystal potential

$$\frac{1}{m_{eff}} = \frac{1}{\hbar^2} \frac{d^2E}{dk^2}$$

Notation  
 $m_{eff} = m^*$

Rigorous form  
 $\vec{p} = \hbar \vec{k}$   
 $\vec{F} = \hbar \frac{d\vec{k}}{dt}$   
 + wave propagation

heavy mass  
 light mass

Semiconductors:  $m_{eff} \approx 0.01 - 0.1 m_{electron}$   
 (near bandgap)

- Newton's 2<sup>nd</sup> law not violated for crystal as a whole
- It means when electron has force on it by eg. applied E or B field, it gets accelerated more than you would expect because forces also applied to it by the ions (nuclei) that are causing the overall periodic potential.

tensor

If curvature is different in different directions, they need tensors

$$\left(\frac{1}{m}\right)_{xx} = \frac{1}{\hbar^2} \frac{d^2E}{dk_x^2}$$

$$\left(\frac{1}{m}\right)_{xy} = \frac{1}{\hbar^2} \frac{d^2E}{dk_x dk_y}$$

etc.

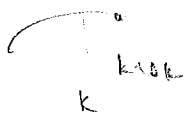
Then N2 is  $F = m \frac{dv}{dt} \rightarrow \frac{dv}{dt} = \frac{1}{m} F \rightarrow \frac{dv_x}{dt} = \left(\frac{1}{m}\right)_{xx} F_x + \left(\frac{1}{m}\right)_{xy} F_y + \left(\frac{1}{m}\right)_{xz} F_z$

$\uparrow$                        $\uparrow$                        $\uparrow$   
 x                                  y                                  z  
 etc.

$m_{eff}$  can even be negative!



what it means = going from  $k$  to  $k+2k$



the lattice causes a bigger  $\frac{d^2E}{dk^2}$  than the applied force does

Convert old eqns

$$\sigma = \frac{ne^2\tau}{m} \rightarrow \sigma = \frac{nq^2\tau}{m^*}$$

small  $m \rightarrow$  better conductor

$$\omega_c = \frac{eB}{mc} \rightarrow \omega_c = \frac{eB}{m^*c}$$

one way to measure  $m^*$  is cyclotron resonance

$$D(E) = \frac{V}{2\pi^2} \left( \frac{2m^*}{\hbar^2} \right)^{3/2} \sqrt{E}$$

$$E_F = \left( \frac{3m^2N}{V} \right)^{2/3} \left( \frac{\hbar^2}{2m^*} \right) \quad \times$$

etc.

Holes

like auditorium w/ one empty seat, (also like bubbles.)



electron picture  
 $m_e = \text{negative}$

$\Rightarrow$  hole picture



$$E_h = \frac{\hbar^2 k_h^2}{2m_h}$$

$m_h = \text{positive}$

$$\vec{k}_h = -\vec{k}_e$$

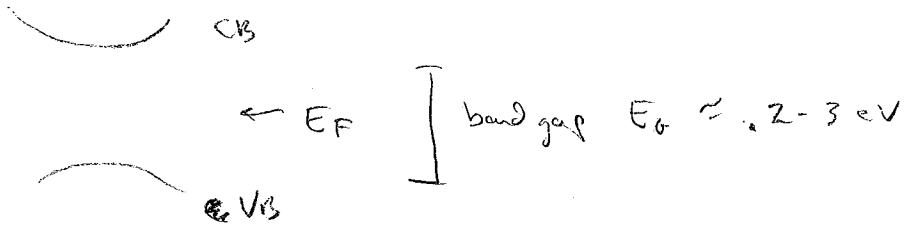
$E_h = -E_e \rightarrow$  band flipped upside down

charge =  $+|e|$

spin = also reversed

people frequently talk of holes, but don't usually draw bands like  $\cup$  valence.

Valence + Conduction bands



General Rule (Problem 9.3) -  $m^* \sim E_g$ , approximately for direct gap materials  
 ↓  
 1) may or may not assign, uses 2<sup>nd</sup> order perturbation theory, uses perfect relations, would need to translate

Table of  $E_g$  = pg 190 Table 1

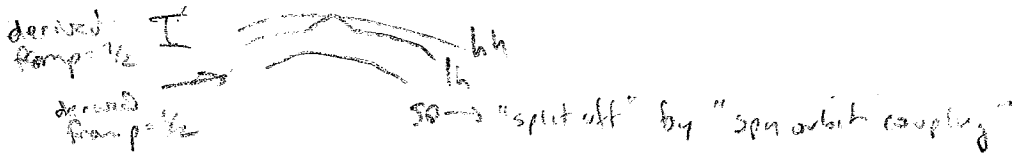
Table of  $m^*$  = pg 201 Table 2

Handout of  
 Direct gap semiconductors?  
 Figure? Section?

VB details

See Fig 14 pg 205 for Ge

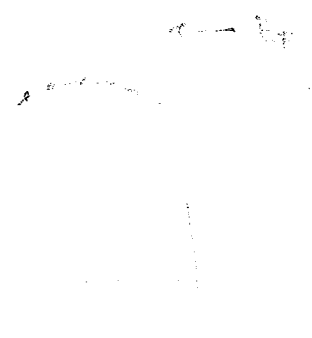
or needed for Si + GaAs?



~~Handwritten scribbles~~

day 33 pg 4

← how many electrons available for con?



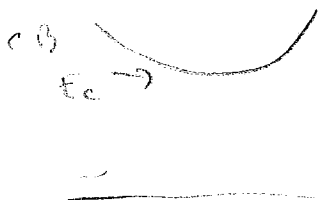
$$f = \frac{1}{e^{-E_c/kT} + 1}$$

"temp of ground"  $E_c \gg kT$

so already almost all electrons are in ground

$$f \approx \frac{1}{e^{E_c/kT}}$$

Side note: maybe after next class



← density of states

just like free electron except

→  $2m^*$  to explain difference in carrier concentration

→  $E - E_c$  to explain vertical offset

→ done in Chap 6 hand out

$$D(E) = \frac{1}{2\pi^2} \left( \frac{2m^*}{\hbar^2} \right)^{3/2} (E - E_c)^{1/2}$$

concentration  $n = \int_{E_c}^{\infty} \left( \frac{D(E)}{V} \right) f(E) dE$

$$= \int_{E_c}^{\infty} \left( \frac{1}{2\pi^2} \right) \left( \frac{2m^*}{\hbar^2} \right)^{3/2} (E - E_c)^{1/2} e^{-(E - E_c)/kT} dE$$

$$= \frac{1}{2\pi^2} \left( \frac{2m^*}{\hbar^2} \right)^{3/2} e^{+E_c/kT} \int_{0}^{\infty} (E - E_c)^{1/2} e^{-E/kT} dE$$

Mathematical  $\frac{\sqrt{\pi}}{2} e^{-E_c/kT} (kT)^{3/2}$

Eqn 8.39 pg 206

$$n = 2 \left( \frac{m^* kT}{2\pi\hbar^2} \right)^{3/2} e^{-(E_c - \mu)/kT}$$

$$= n_0 e^{-(E_c - \mu)/kT}$$

$$n_0 = 2 \left( \frac{m^* kT}{2\pi\hbar^2} \right)^{3/2}$$

pg 214

proof of 2.5.

$$\frac{2 \cdot 2^{3/2}}{2^{-1/2}}$$

$$\frac{2^{3/2}}{2} = 2^{1/2}$$

$$\frac{1}{2} = 2^{-1/2}$$

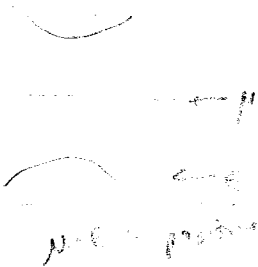
$$= 2^{1/2} \checkmark$$

problem: don't know p.

Do some thing for holes in valence band

$$f_v(E) = 1 - f_c(E)$$

$$= 1 - \frac{1}{e^{(E_c - E)/kT} + 1} = \frac{e^{(E - E_c)/kT}}{e^{(E - E_c)/kT} + 1}$$



(E = energy of electrons, μ = chemical energy of holes which must have opposite sign)

$$\frac{e^x}{e^x + 1} = \frac{e^{-x}}{e^{-x} + 1} = \frac{1}{1 + e^x}$$

$$= \frac{1}{e^{(E_c - E)/kT} + 1}$$

$$= e^{-(E - E_c)/kT}$$

Same integral Type pg 206

$$p = \int_{-\infty}^{E_v} \frac{D_v(E) f_v(E)}{V} dE$$

mod for valence

pg 206  
Eqn 8.42  
Typical Kittel!

$$p = 2 \left( \frac{m_{mv}^* kT}{2\pi \hbar^2} \right)^{3/2} e^{-(E_v - E_f)/kT}$$

(Same integral)

could be called p<sub>0</sub>, but isn't in Kittel. I'll use that symbol later, though

Multiply together

day 34  
pg 1

$$np = 4 \left( \frac{kT}{2\pi \hbar^2} \right)^3 \left( \frac{m_e m_h}{m_0^2} \right)^{3/2} e^{-E_g/kT}$$

$$np = 4 \left( \frac{kT}{2\pi \hbar^2} \right)^3 (m_e m_h)^{3/2} e^{-E_g/kT}$$

n<sub>0</sub>p<sub>0</sub>

only assumption: μ = E<sub>F</sub> for both CB & VB

could be "intrinsic", could be "degenerate"

↳ but conditions 1<sup>st</sup> condition is full wet "uncompensated"